Impact of near-inertial waves on vertical mixing and air-sea CO₂ fluxes in the Southern Ocean

Hajoon Song¹, John Marshall², Jean-Michel Campin² and Dennis J. McGillicuddy Jr.³

3

9

Key Points:

4	¹ Department of Atmospheric Sciences, Yonsei University, Seoul, Republic of Korea
5	² Department of Earth, Atmospheric and Planetary Sciences, Massachusetts Institute of Technology, Cambridge,
6	Massachusetts, USA
7	³ Department of Applied Ocean Physics and Engineering, Woods Hole Oceanographic Institution, Woods Hole,
8	Massachusetts, USA

Filtering out near-inertial waves from a simulation diminishes vertical shear, shallowing the mixed layer depth. The shallowed mixed-layer reduces the supply of dissolved inorganic carbon to the surface. These high frequency processes can change CO₂ flux by 1/3 amplitude of the sea-sonal cycle.

Corresponding author: H. Song, hajsong@yonsei.ac.kr

16 Abstract

We report the significant impact of near-inertial waves (NIWs) on vertical mixing and 17 air-sea carbon dioxide (CO₂) fluxes in the Southern Ocean (SO) using a biogeochemical 18 model coupled to an eddy-rich ocean circulation model. The effects of high-frequency 19 processes are quantified by comparing the fully coupled solution (ONLINE) to two offline 20 simulations based on five-day averaged output of the ONLINE simulation: one that uses 21 vertical mixing archived from the ONLINE model (CTRL) and another in which verti-22 cal mixing is recomputed from the five-day average hydrodynamic fields (5dAVG). In this 23 latter simulation, processes with temporal variabilities of a few days including NIWs are 24 excluded in the biogeochemical simulation. Suppression of these processes reduces verti-25 cal shear and vertical mixing in the upper ocean, leading to decreased supply of carbon-26 rich water from below, less CO₂ outgassing in austral winter, and more uptake in summer. 27 The net change amounts up to 1/3 of the seasonal variability in SO CO₂ flux. Our results 28 clearly demonstrate the importance of resolving high frequency processes such as NIWs to 29 better estimate the carbon cycle in numerical model simulations. 30

1 Introduction

Intensive vertical mixing is one of the important aspects that characterize the South-32 ern Ocean (SO) [de Boyer Montégut et al., 2004; Holte et al., 2017]. To the north of the 33 Antarctic Circumpolar Current (ACC), thick mixed layers in austral winter reach a few 34 hundred meters deep and contribute to the exchanges of momentum, heat and freshwater, 35 and formation of Antarctic Intermediate Water and Subantarctic Mode Water [Dong et al., 36 2008]. Deep vertical mixing is also important in biogeochemical processes. It brings sur-37 face water, rich in oxygen, to the interior ocean during the process of intermediate and 38 mode water formations and forms the oxygen maximum layer [Talley et al., 2011]. The 39 uptake of anthropogenic tracers (e.g. chlorofluorocarbons) can be expedited through deep 40 convection [Shao et al., 2013; Song et al., 2015]. Intensive vertical mixing is also efficient 41 in drawing the subsurface water mass rich in nutrients and carbon close to the surface. 42 For example, iron, critical for the primary production in the SO, is supplied from the in-43 terior of the ocean by vertical mixing in winter [Tagliabue et al., 2014]. Vertical mixing 44 is claimed to be responsible for the nitrate transport from the deep Eastern South Pacific 45 to the Patagonian shelf region that is one of the most productive areas in the world [Song 46 et al., 2016a]. 47

-2-

48	Near-inertial waves (NIWs) can significantly impact surface vertical mixing. NIWs
49	have a frequency close to an inertial frequency with a length scale of 10 to 100 km and
50	are mainly excited by wind forcing [Alford et al., 2016]. In particular, midlatitude storms
51	can drive the ocean with their inertially rotating components [D'Asaro, 1985] and excite
52	high and low mode waves. High modes create counterclockwise rotation in the SO, pro-
53	ducing high vertical shear of currents and enhanced vertical mixing with the downward
54	propagation of energy [Alford et al., 2016; Alford and Gregg, 2001]. In numerical exper-
55	iments, NIWs deepen mixed layer depths (MLDs) by up to 30% on annual average be-
56	tween 40°S and 60°S [Jochum et al., 2013]. This deepening of MLDs in the SO is asso-
57	ciated with the midlatitude storm tracks [Simmons and Alford, 2012], which is consistent
58	with the view that a high level of wind work can generate NIWs.

The modulation of vertical mixing by NIWs are expected to influence the air-sea 59 CO_2 exchange in the SO. The SO CO_2 flux is controlled by two key processes: biologi-60 cal drawdown and intensive vertical mixing [Wetzel et al., 2005; Takahashi et al., 2009]. 61 In austral summer, the SO takes up CO₂ from the atmosphere with biological processes 62 leading to drawdown, but it emits CO2 into the atmosphere in austral winter due to the 63 entrainment of carbon-rich interior water up to the surface. It is hence anticipated that en-64 hanced vertical mixing by NIWs promotes more entrainment in austral winter, resulting 65 in more CO₂ outgassing. However, it is not well investigated how much these processes 66 influence air-sea CO2 exchange. In addition, their role in austral summer is uncertain as 67 biological drawdown is itself impacted by the MLD. Deepening of the MLD in austral 68 summer may enrich the surface ocean with carbon, but at the same time, bring more nu-69 trients that potentially promote higher biological drawdown. If we consider the changes in 70 temperature by mixing that regulates the solubility, it becomes more difficult to anticipate 71 the role of high frequency processes on CO₂ flux. 72

Here we attempt to quantify the integral effect of NIWs on CO₂ flux near the Drake 73 Passage. Our approach is to compare CO₂ fluxes with and without high temporal vari-74 abilities including NIWs by making use of both online and offline biogeochemical mod-75 els. The online biogeochemical simulation is forced by 6-hourly wind product to properly 76 simulate the input of energy with near inertial frequency. The offline biogeochemical sim-77 ulation is forced by the same wind forcing, but the ocean states driving biogeochemical 78 variables are the 5-day averaged fields from the online simulation, as described in section 79 2. The results in section 3 suggest that they are capable of altering the air-sea CO_2 flux 80

-3-

significantly through the vertical flux of carbon. Then the discussion on air-sea CO_2 simulation follows in section 4.

⁸³ 2 Simulations of air-sea CO₂ flux

84

2.1 Eddy-Resolving Model

A simple biogeochemical model [Dutkiewicz et al., 2005; Parekh et al., 2006; Verdy 92 et al., 2007] is coupled to a 1/20° version of the MIT Ocean General Circulation Model 93 (MITgcm) [Marshall et al., 1997a,b; Adcroft et al., 1997; Marshall et al., 1998; Adcroft 94 et al., 2004]. The model domain covers the area from -75°S to -35°S with a 140° longitu-95 dinal swath of a section of the Antarctic Circumpolar Current (ACC) stretching from the 96 Southeast Pacific through Drake Passage to the Southwest Atlantic (Fig. 1(a)) and has 50 97 vertical levels with higher resolution near the surface (10 m). The model configuration is 98 that described in Tulloch et al. [2014] and is same as that used in Song et al. [2016b]. In 99 particular, the ocean model was integrated with 6-hourly wind and buoyancy fluxes from 100 the reanalysis data set by European Centre for Medium-Range Weather Forecasts [Sim-101 mons et al., 2007]. Vertical mixing is parameterized by the nonlocal K-profile parameter-102 ization (KPP) scheme of Large et al. [1994] in which the depth of the boundary layer is 103 first estimated based on the bulk Richardson number determined by the surface forcing, 104 buoyancy and vertical shear of velocity. Then the vertical diffusivity is calculated within 105 the boundary layer proportionally to its depth. The biogeochemical boundary conditions 106 are provided from the monthly mean states from a global model [Song et al., 2016b]. The 107 integration of this configuration (referred to as ONLINE hereinafter) for 4 years with a 108 time step of 2 minutes results in the ocean states that have similarities to the observed SO 109 states including the high level of mesoscale eddy activities whose spatial scale is O(100)110 km). 111

In the rotary spectra of velocity shear at 100 m in the model, there are two distinctive lines with elevated power at each latitude: one for the frequency band near zero and the other close to the inertial frequency (Fig. 2a). While the former is associated with geostrophic component, the latter reveals NIWs. Velocity shear rotates counterclockwise with positive frequencies near inertial frequency, f. The rotary spectra of velocity also show enhanced variance near f as well as low frequency (not shown), consistent with that obtained from drifter velocity observations [*Elipot and Lumpkin*, 2008]. The wavenum-

-4-



(b) Dissolved inorganic carbon, September



Figure 1. (a) A snapshot of sea surface temperature from the online simulation (ONLINE) in September. The mask in the gray color scale around Antarctica represents the sea ice concentration. The rectangular box at the center indicates the region for the vertical section of dissolved inorganic carbon from the surface to 1500 m in (b). The thin black line in (b) marks the mixed layer depth (MLD) defined as the level whose density is 0.03 kg m^{-3} greater than the surface value. Above the vertical section is the CO₂ flux and the mask in the gray color scale at 0 m is the sea ice concentration. The direction of the ACC is shown as a gray arrow in (b).

- ber spectrum using the surface current shows the slope close to k^{-3} , with k being the
- ¹²⁰ wavenumber (Fig. 2(c)). Having this slope in the spectrum suggests that mesoscale eddies

derive their energy source from baroclinic instability [*Sasaki et al.*, 2014].

127

2.2 Simulation of air-sea CO₂ flux

128

The air-sea CO₂ flux (F_{CO_2}) is estimated using a gas exchange parameterization:

 $F_{\rm CO_2} = K_w (1 - A_{SI}) \left(p {\rm CO}_2 - p {\rm CO}_2^{atm} \right), \tag{1}$

where K_w is the gas transfer velocity (m s⁻¹) determined by wind speed squared and sea surface temperature [*Wanninkhof*, 1992], A_{SI} is the fraction of the sea ice coverage within a model grid cell varying from 0 when there is no sea ice to 1 when the grid cell is fully covered by sea ice, and pCO_2 and pCO_2^{atm} are the oceanic and atmospheric partial pressure of CO₂, respectively. We fix pCO_2^{atm} at the pre-industrial level (278 ppm) and oceanic surface pCO_2 is estimated using dissolved inorganic carbon (DIC), alkalinity as well as temperature and salinity, following *Follows et al.* [2006].

The distribution of DIC in austral winter reveals the upwelling that increases the 137 surface DIC poleward (Fig. 1(b)). Near Antarctica where the surface DIC is the great-138 est, the CO_2 outgassing is inhibited by the presence of sea ice. Instead, the strongest CO_2 139 outgassing occurs near the ACC where the intense vertical mixing increases the surface 140 DIC and, hence, pCO_2 . Strong westerly wind forcing also contributes to effective CO_2 141 outgassing adjacent to the ACC. Interestingly, the sign of CO2 flux changes near 55°S 142 where the MLD abruptly shallows. The region north of that latitude takes up CO_2 from 143 the atmosphere as pCO_2 becomes smaller than pCO_2^{atm} . Deep mixing can provide iron 144 [Tagliabue et al., 2014] and promote biological drawdown that partially compensates the 145 increase of the surface pCO_2 , but its magnitude in our simulation is much less than that 146 of DIC vertical flux because of light limitation in winter. Mesoscale eddies in the model 147 simulations alter DIC concentration as shown in the fluctuating isolines of DIC and hence 148 modulate CO_2 flux by perturbing the concentration of DIC. See Song et al. [2016b] for a 149 detailed discussion of the modulation of CO_2 flux by the mesoscale. 150

151

2.3 Suppression of NIWs in the biogeochemical model

152

153

For the quantification of the impact by NIWs on air-sea CO_2 flux, we designed an offline simulation of the biogeochemical model with suppressed variances in time. This



Figure 2. The rotary spectra of vertical shear ($s^{-2} \text{ cpd}^{-1}$) in \log_{10} scale at 100 m in the (a)1/20° eddyresolving model and (b) 5-day averaged velocity. Blue lines in (a, b) represent the inertial frequency, $-f/(2\pi)$, where *f* is the Coriolis frequency. The kinetic energy spectra of the surface current in ONLINE (red) and 5dAVG (blue) are plotted in (c). Shading indicates the maximum/minimum energy levels during the last 3-year simulation. The inset in (c) is the surface vorticity from 5dAVG.

offline simulation (referred to as 5dAVG hereinafter) was integrated from the second year of ONLINE for 3 years driven by 5-day averaged temperature (T), salinity (S) and horizontal velocities (U, V) from ONLINE. The vertical mixing for tracers was recalculated by the KPP scheme using those 5-day averaged ocean states. The surface forcing is not changed in the surface mixing model and the biogeochemical model.

The 5-day average significantly suppresses the processes with the time scale shorter 159 than a few days, and the peak along f in the variances disappears (Fig. 2(b)). Instead al-160 most all the energy in the vertical shear in 5dAVG is concentrated near zero frequency 161 related to the geostrophic component. The 5-day average slightly lowers the kinetic energy 162 in the wavenumber band corresponding to the mesoscale as well (Fig. 2(c)). However the 163 surface vorticity in 5dAVG shown in the inset of Figure 2(c) shows rich variabilities asso-164 ciated with the mesoscale. Hence, we can interpret the solution from this offline biogeo-165 chemical model having a similar phenomenology as ONLINE except the surface vertical 166 mixing with the lack of the energy associated with NIWs. 167

As the focus of this study lies on the impact of the vertical mixing by NIWs on the 168 CO2 flux, we designed additional offline biogeochemical simulation where the surface ver-169 tical mixing (eddy diffusivity and eddy diffusivity weighted by non-local transport coef-170 ficient) from ONLINE were provided to the offline simulation. This offline simulation is 171 referred to as CTRL hereinafter. This CTRL simulation then has the same mean physical 172 oceanic states and surface vertical mixing as ONLINE, and the comparison of CTRL and 173 5dAVG allows us to isolate the impact of vertical mixing associates with NIWs on CO2 174 flux. All three biogeochemical simulations are summarized in Table 1. 175

179

3 Impact of NIWs on Mixed Layer Depth and air-sea CO₂ fluxes

180

3.1 Mixed Layer Depth

The CTRL run captures the observed wintertime MLD structure: deep MLDs reaching deeper than 500 m upstream of Drake Passage (Fig. 1(b)) and relatively shallow MLDs of approximately 150 m in the Atlantic (Fig. 3(a)) [*de Boyer Montégut et al.*, 2004; *Dong et al.*, 2008; *Holte et al.*, 2017]. The absence of NIWs in 5dAVG weakens the vertical mixing and makes the MLD shallower. The reduction in MLD occurs almost everywhere, but it is more pronounced along the ACC where it can be greater than 100 m (Fig. 3(b)).

-8-



Figure 3. (a) Mixed layer depth (MLD) in the CTRL simulation averaged in September and (b) the changes in MLD between CTRL and 5dAVG. The mask in gray scale is the simulated sea ice fraction. (c) and (d) are histograms of MLDs in CTRL (red) and 5dAVG (blue) in January and September, respectively, with the mean MLD differences (black bar plots) between the two runs binned by the original MLDs in CTRL. Black lines in the bar plots represent the standard errors. Note that the scales in x and y axes are different between (c) and (d).

176 **Table 1.** Comparison of biogeochemical simulations. In the table, T, S, u and v represent a snapshot of

temperature, salinity, zonal velocity and meridional velocity, respectively, with a frequency of 120 seconds.

The 5-day average of those variables are written using angled bracket, $\langle \rangle$.

Name	ONLINE	CTRL	5dAVG
Coupling method to the circula-	online	offline	offline
tion model			
Physical fields for biogeochem-	T, S, u, v	$\langle T\rangle,\langle S\rangle,\langle u\rangle,\langle v\rangle$	$\langle T\rangle, \langle S\rangle, \langle u\rangle, \langle v\rangle$
istry			
Eddy diffusivity (κ) in the KPP	computed using	$\langle \kappa \rangle$ from ON-	computed using
scheme	T, S, u and v	LINE is loaded.	$\langle T \rangle, \langle S \rangle, \langle u \rangle$ and
			$\langle v \rangle$

The reduction in MLD in 5dAVG can be further evaluated using the histograms of 193 MLDs. In January, the mean MLD in 5dAVG is 12.4 m, approximately 5 m shallower 194 than that in CTRL (Fig. 3(c)). The histogram of 5dAVG is shifted to the left and be-195 comes more skewed: the moment coefficient of skewness is increased from 1.0 to 1.5 196 when the vertical mixing is recalculated using 5-day averaged fields. The size of MLD 197 reduction increases with the background MLD in CTRL until 70 m, but decreases be-198 yond that point (the black bar plot in Fig. 3(c)). Similar patterns stand out in the MLD 199 changes in September (Fig. 3(d)). The mean MLD in 5dAVG is 83.5 m, which is roughly 200 24 m shallower than that in CTRL. The histogram of MLD in 5dAVG is more positively 201 skewed than that in CTRL: the moment of coefficient of skewness in 5dAVG is 16.6 while 202 it is 14.7 in CTRL. The reduction of MLD tends to increase as the background MLD be-203 comes deeper, but it starts to decrease with the background MLD for the bins greater than 204 roughly 300 m. 205

We argue that the reduction of MLD results from the elimination of vertical shear of the current in the near inertial frequency band associated with NIWs (Fig. 2a,b). The depth of surface boundary layer in the KPP mixing scheme is defined as the level where the Richardson number is smaller than a predefined critical value. Since the Richardson number is inversely proportional to the vertical shear of the current, the decrease of verti-

cal shear in 5dAVG can result in the increase of the Richardson number and the reduction 211 of the MLD. Just as the temporal smoothing suppresses the kinetic energy, especially in 212 the mesoscale (Fig. 2(b)), it can alter the stratification near the surface. Suppose, for ex-213 ample, there is a cold eddy that moves along the current in the region of a constant strat-214 ification. Taking a 5-day average of the temperature has the effect of the introduction of 215 cold (and heavy) water to the path over the 5-day span and can change the stratification 216 and hence, MLD. However, we find that the 5-day average has little impact on squared 217 Brunt-Väisälä frequency (N²) in the upper ocean (not shown), suggesting that the MLD 218 reduction is driven mainly by the reduced level of vertical shear with the absence of NIWs 219 in 5dAVG. 220

221

3.2 Air-sea CO₂ flux

The CO_2 flux in ONLINE exhibits seasonal variability (red line in Fig. 4(a)). In 222 summer, the ocean takes up CO_2 from the atmosphere while the ocean emits CO_2 back 223 to the atmosphere in winter. This is consistent with previous findings on air-sea CO2 ex-224 change in the SO [Lenton et al., 2013]. In CTRL where the 5-day averaged fields includ-225 ing vertical mixing drive the biogeochemical model, the ocean takes up more CO2 in sum-226 mer and releases more CO_2 in winter (green line in Fig. 4(a)). Although the change in 227 CO_2 flux can be as big as 0.1 Pg C yr⁻¹, the mean difference is rather small (-0.04 Pg 228 C yr⁻¹), suggesting that CTRL is a realistic approximate to ONLINE. When the surface 229 vertical mixing is recalculated using the 5-day averaged fields in 5dAVG, the CO₂ uptake 230 in summer is further enhanced while the outgassing of CO_2 in winter weakens (blue line 231 in Fig. 4(a)). The mean change in CO₂ flux in 5dAVG is -0.23 Pg C yr⁻¹ with respect to 232 ONLINE and -0.19 Pg C yr⁻¹ with respect to CTRL, with the maximum reduction oc-233 curring in winter. The size of the CO2 flux change in 5dAVG is considerable, with an am-234 plitude of 30% of the seasonal cycle of CO₂ in ONLINE. It is also large when compared 235 with the seasonal cycle of CO₂ in CTRL, roughly 20% of that level. 236

Suppressing processes with high temporal frequency is responsible for the shift of the CO₂ flux curve of 5dAVG downward compared to that of ONLINE (Fig. 4(a)). Because the surface wind is fixed and the mean temperature is the same, the downward shift of the CO₂ flux curve of 5dAVG is solely the result of the reduction in the surface pCO_2 . The impact of removing NIWs alone is quantified by comparing CTRL and 5dAVG simulations. Bigger impact is observed in winter and spring, while the summertime shows the



Figure 4. An area integrated (a) CO₂ flux and (b) diapycnal mixing of DIC for each month from ONLINE (red), CTRL (green) and 5dAVG (blue).

minimum impact. This seasonality in the impact of suppressed NIWs is to be expected,
given that the main energy source for NIWs is the wind.

To quantitatively investigate the driver of the CO₂ flux changes, we write the DIC concentration thus:

249

$$\frac{\partial \text{DIC}}{\partial t} = -\nabla \cdot (\mathbf{u}\text{DIC}) + \kappa \frac{\partial^2 \text{DIC}}{\partial z^2} - F_{\text{CO}_2} + S_{\text{bio}} + S_{\text{C}}, \qquad (2)$$

where **u** is a 3-dimensional velocity vector, κ is the vertical diffusivity, F_{CO_2} is the CO₂ 250 flux, S_{bio} and S_C are the source/sink terms associated with biological activities and cal-251 cium carbonate flux, respectively. According to (2), the tendency of DIC is determined by 252 advection, diapycnal mixing and addition/subtraction of DIC through biogeochemical pro-253 cesses and air-sea exchange whose rate is determined by the current pCO_2 level through 254 (1). When comparing the terms in ONLINE, CTRL and 5dAVG, excluding F_{CO_2} , the 255 biggest differences are found in $\kappa \partial^2 \text{DIC}/\partial z^2$ (Fig. 4(b); other terms not shown). Clearly 256 the magnitude of the changes in diapycnal mixing term is sufficient to explain the CO₂ 257 flux changes, indicating that the supply of carbon-rich water from below through vertical 258 mixing is the primary driver. 259

The changes in CO_2 flux is the largest in the vicinity of the ACC. In summer when almost all regions in the model domain take up CO_2 , suppressing processes with high temporal variability allows the ocean to absorb more CO_2 near the ACC (Fig. 5(a,c)). Although the changes in CO_2 flux north of the ACC and Brazil-Malvinas confluence zone are positive, they are not big enough to convert the CO_2 flux from the uptake to outgassing in CTRL. When NIWs are suppressed, the uptake of CO_2 becomes even stronger



Figure 5. (a,b) Monthly averaged CO_2 flux in ONLINE, and (c,d) the difference between CTRL and ON-LINE, and (e,f) between 5dAVG and CTRL. The panels on the left (a,c,e) are the CO_2 flux and its difference in January, and the panels on the right (b,d,f) are those in September. Similar to Figure 1, the gray mask near Antarctica represents the sea ice concentration.

mainly near the ACC in summer (Fig. 5(e)). In winter, using 5-day averaged physical
fields and eddy diffusivity in the biogeochemical model does not alter the CO₂ flux as
much as in summer near the ACC (Fig. 5(b,d)). The biggest changes are in the BrazilMalvinas confluence zone. There is originally CO₂ uptake in ONLINE (Fig. 5(b)), but
the positive difference between CTRL and ONLINE suggests CO₂ flux is close to zero
or even positive (outgassing) in CTRL (Fig. 5(d)). The differences between 5dAVG and
CTRL are negative (Fig. 5(f)), indicating that suppressing NIWs reduces CO₂ outgassing.

These spatial changes are closely linked to the changes in the vertical diffusive flux 277 of carbon by diapycnal mixing. In summer, although the biological pump is the dominant 278 term in (2) and leads to CO_2 uptake, the vertical diffusive flux tends to increased DIC at 279 the surface (Fig. 6(a)). The difference in vertical diffusive flux of carbon between CTRL 280 and ONLINE is generally negative, indicating the supply of carbon through vertical mix-281 ing is even weaker when high temporal variability is suppressed. 5dAVG shows even more 282 reduction of carbon supply by vertical mixing when filtering out NIWs (Fig. 6(e)). The 283 responses of the biological pump to suppressing high temporal variability and NIWs are 284 not as large as those of vertical mixing (not shown), hence reduced supply of carbon from 285 below leads to more uptake of CO_2 (Fig. 4(a) and 5(c,e)). 286

Changes in vertical diffusive flux of carbon can also explain the CO₂ flux changes 287 between simulations in winter. In ONLINE, the carbon supply by diapycnal vertical mix-288 ing is the dominant signal at the upstream of the Drake Passage (Fig. 6(b)). When the 289 biogeochemical model is driven by physical fields with the same mean states (including 290 vertical mixing) but without high temporal variabilities (CTRL), the carbon supply to the 291 surface by vertical mixing is enhanced in the upstream of the Drake Passage and Brazil-292 Malvinas confluence zone (Fig. 6(d)). However, when NIWs are suppressed (5dAVG), the 293 vertical diffusive flux of carbon is reduced compared with CTRL (Fig. 6(f)) and even ON-294 LINE (not shown). The reduction of carbon flux is particularly pronounced near the Drake 295 Passage and vicinity of the ACC. 296

²⁹⁸ The results from ONLINE and 5dAVG demonstrate that suppressing NIWs and pro-²⁹⁹ cesses of high temporal variability shallows the MLD as shown in Figure 3. Since eddy ³⁰⁰ diffusivity, κ , in the KPP scheme is proportional to the depth of surface boundary layer, ³⁰¹ it is the smallest in 5dAVG, leading to the weakest vertical diffusive flux of carbon to the ³⁰² surface. The decrease of carbon supply from subsurface lowers the *p*CO₂ at the surface



Figure 6. Same as Figure 5, but for the vertical mixing contribution to the surface DIC.

297

303	and causes more uptake of CO ₂ in summer and less CO ₂ outgassing in winter. Driving
304	the biogeochemical model with the same mean states including vertical mixing but with-
305	out processes of high temporal variability (CTRL) also modifies the diffusive flux of car-
306	bon, lowering surface pCO_2 in summer but increasing it in winter (Fig. 4(b), 6(c,d)). The
307	effect of NIWs on CO_2 flux can be isolated by comparing CTRL and 5dAVG, and the
308	comparison suggests that the effect of NIWs in CO_2 flux is greater in winter (Fig. 4(b),
309	6(e,f)). The decrease of carbon supply by vertical mixing can explain the reduced out-
310	gassing of CO_2 when NIWs are suppressed in winter (Fig. 5(f)).

4 Discussion

Three numerical experiments demonstrate the sensitivity of air-sea CO_2 flux to highfrequency processes. Weakening of vertical mixing leads to less supply of carbon-rich water from below and reduces the surface carbon concentration. Budget analysis of dissolved inorganic carbon (DIC) reveals that the diffusive carbon flux by vertical mixing is systematically lower when NIWs are suppressed (5dAVG; Fig. 4b). As a result, the SO emits less CO_2 in austral winter and takes up more CO_2 in austral summer, and this change is approximately 1/3 of seasonal variability in ONLINE.

We note that removing NIWs and processes with high temporal variability can po-319 tentially affect the air-sea CO₂ exchange by changing the solubility of the surface ocean. 320 Temporal smoothing reduces the warm/cold temperature anomalies and increases/decreases 321 solubility, resulting in the decrease/increase of partial pressure of CO₂ that changes CO₂ 322 flux. However, we expect a very small net effect. The temporal smoothing preserves mean 323 properties, hence there is no net sea surface temperature change. Since the changes in par-324 tial pressure of CO₂ with respect to temperature is fairly constant [Takahashi et al., 1993], 325 net-zero temperature change should yield a very small net change in partial pressure of 326 CO_2 . 327

We show that processes with a high temporal frequency such as NIWs have a significant impact on vertical mixing and CO₂ flux, suggesting that these processes must be resolved in numerical simulations for better estimation of CO₂ flux in the SO. In order to properly simulate NIWs, the frequency of the wind forcing should be higher than the local inertial frequency. The inertial frequency increases with latitude, and it is close to 2 cycles per day near -75° S, the southern boundary of our model (Fig. 2a). At this latitude, at least 6-hourly winds are required to accurately force NIWs. However, many models do not use wind forcing with the frequency adequate for resolving NIWs, including those in Coupled Model Intercomparison Project phase 5 (CMIP5): almost half of coupled climate models have daily coupling frequency [*Tian*, 2016]. Even eddy-resolving models may miss the energy associated with NIWs if the wind forcing frequency is lower than local inertial frequencies.

Our study demonstrates that NIWs have broad impacts not only on the surface ocean states but on the carbon cycle and air-sea CO₂ exchange through vertical mixing. The impact is the strongest near the ACC in winter where many climate models show a shallow MLD bias [*Downes et al.*, 2015]. The underestimation of the MLD can lead to not enough uptake of transient tracers (e.g., chlorofluorocarbons), and our study suggests that this could be in part due to the inability to fully resolve NIWs and associated mixing.

In recent years, the Southern Ocean Carbon and Climate Observations and Modeling (SOCCOM) project deployed floats with the ability to estimate the carbon concentration at the surface and CO₂ exchange between the atmosphere and the ocean. The new observations revealed that the wintertime CO₂ outgassing near the ACC had been underestimated by most numerical models. Hence it is important to evaluate whether NIWs are properly included in those numerical simulations in diagnosing the models' performance.

352 Acknowledgments

The MITgcm can be obtained from http://mitgcm.org. The surface forcing data used in 353 this study are available at the Geophysical Fluid Dynamics Laboratory data portal web-354 page (http://data1.gfdl.noaa.gov/nomads/forms/core/COREv2.html). Resources support-355 ing this work were provided by the NASA High-End Computing (HEC) Program through 356 the NASA Advanced Supercomputing (NAS) Division at Ames Research Center with 357 the award number SMD-15-5752. HS, JM, and DJM were supported by the NSF MOBY 358 project (OCE-1048926 and OCE-1048897). HS acknowledges the support by Yonsei Uni-359 versity Research Fund of 2018-22-0053. DJM also gratefully acknowledges NASA sup-360 port. 361

362 References	
-----------------------	--

363	Adcroft, A., C. Hill, and J. Marshall (1997), Representation of topography by shaved cells
364	in a height coordinate ocean model, Mon. Wea, Rev., 125, 2293–2315.

- Adcroft, A., C. Hill, J.-M. Campin, J. Marshall, and P. Heimbach (2004), Overview of the
- formulation and numerics of the MIT GCM, in *Proceedings of the ECMWF Seminar Se*
- ries on Numerical Methods, Recent Developments in Numerical Methods for Atmosphere
- and Ocean Modelling, pp. 139–149, ECMWF.
- Alford, M. H., and M. C. Gregg (2001), Near-inertial mixing: Modulation of shear, strain and microstructure at low latitude, *J. Geophys. Res. Oceans*, *106*(C8), 16,947–16,968,
- doi:10.1029/2000JC000370.
- Alford, M. H., J. A. MacKinnon, H. L. Simmons, and J. D. Nash (2016), Near-Inertial
- Internal Gravity Waves in the Ocean, Annu. Rev. Mar. Sci., 8(1), 95–123, doi:
- ³⁷⁴ 10.1146/annurev-marine-010814-015746.
- ³⁷⁵ D'Asaro, E. (1985), The energy flux from the wind to near-inertial motions in the mixed ³⁷⁶ layer, *J. Phys. Oceanogr.*, *15*, 1043–1059.
- de Boyer Montégut, C., G. Madec, A. S. Fischer, A. Lazar, and D. Iudicone (2004), Mixed layer depth over the global ocean: An examination of profile data and a profile-based
- climatology, J. Geophys. Res., 109, C12,003, doi:10.1029/2004JC002378.
- ³⁸⁰ Dong, S., J. Sprintall, S. T. Gille, and L. Talley (2008), Southern Ocean mixed-layer depth from Argo float profiles, *J. Geophys. Res.*, *113*, C06,013.
- Downes, S. M., R. Farneti, P. Uotila, S. M. Griffies, S. J. Marsland, D. Bailey, E. Behrens,
- M. Bentsen, D. Bi, A. Biastoch, C. Böning, A. Bozec, V. M. Canuto, E. Chas-
- signet, G. Danabasoglu, S. Danilov, N. Diansky, H. Drange, P. G. Fogli, A. Gusev,
- A. Howard, M. Ilicak, T. Jung, M. Kelley, W. G. Large, A. Leboissetier, M. Long,
- J. Lu, S. Masina, A. Mishra, A. Navarra, A. J. George Nurser, L. Patara, B. L.
- ³⁸⁷ Samuels, D. Sidorenko, P. Spence, H. Tsujino, Q. Wang, and S. G. Yeager (2015),
- An assessment of Southern Ocean water masses and sea ice during 1988–2007
- in a suite of interannual CORE-II simulations, Ocean Modell., 94, 67–94, doi:
- ³⁹⁰ https://doi.org/10.1016/j.ocemod.2015.07.022.
- ³⁹¹ Dutkiewicz, S., A. Sokolov, J. Scott, and P. Stone (2005), A three-dimensional ocean-
- seaice-carbon cycle model and its coupling to a two-dimensional atmospheric model:
- uses in climate change studies, in *Joint Program on the Sci. Policy Global Change*, chap.
- Rep. 122, MIT, Cambridge, Mass.

395	Elipot, S., and R. Lumpkin (2008), Spectral description of oceanic near-surface variability,
396	Geophys. Res. Lett., 35, L05,606, doi:10.1029/2007GL032874.
397	Follows, M. J., S. Dutkiewicz, and T. Ito (2006), On the solution of the carbonate system
398	in ocean biogeochemistry models, Ocean Modelling, 12, 290-301.
399	Holte, J., L. D. Talley, J. Gilson, and D. Roemmich (2017), An Argo mixed layer climatol-
400	ogy and database, Geophys. Res. Lett., 44, 5618-5626, doi:10.1002/2017GL073426.
401	Jochum, M., B. P. Briegleb, G. Danabasoglu, W. G. Large, N. J. Norton, S. R. Jayne,
402	M. H. Alford, and F. O. Bryan (2013), The impact of oceanic near-inertial waves on
403	climate, J. Climate, 26, 2833–2844.
404	Large, W., J. McWilliams, and S. Doney (1994), Oceanic vertical mixing: A review and a
405	model with nonlocal boundary layer parameterization, Rev. Geophys., 32, 363-403.
406	Lenton, A., B. Tilbrook, R. M. Law, D. Bakker, S. C. Doney, N. Gruber, M. Ishii,
407	M. Hoppema, N. S. Lovenduski, R. J. Matear, B. I. McNeil, N. Metzl, S. E.
408	Mikaloff Fletcher, P. M. S. Monteiro, C. Rödenbeck, C. Sweeney, and T. Takahashi
409	(2013), Sea-air CO ₂ fluxes in the Southern Ocean for the period 1990-2009, Biogeo-
410	sciences, 10(6), 4037-4054, doi:10.5194/bg-10-4037-2013.
411	Marshall, J., C. Hill, L. Perelman, and A. Adcroft (1997a), Hydrostatic, quasi-hydrostatic,
412	and nonhydrostatic ocean modeling, J. Geophysical Res., 102(C3), 5733-5752.
413	Marshall, J., A. Adcroft, C. Hill, L. Perelman, and C. Heisey (1997b), A finite-volume,
414	incompressible Navier Stokes model for studies of the ocean on parallel computers, J .
415	Geophysical Res., 102(C3), 5753–5766.
416	Marshall, J., H. Jones, and C. Hill (1998), Efficient ocean modeling using non-hydrostatic
417	algorithms, J. Mar. Syst., 18, 115-134.
418	Parekh, P., M. J. Follows, S. Dutkiewicz, and T. Ito (2006), Physical and biologi-
419	cal regulation of the soft tissue carbon pump, Paleoceanography, 21, PA3001, doi:
420	10.1029/2005PA001258.
421	Sasaki, H., P. Klein, B. Qiu, and Y. Sasai (2014), Impact of oceanic-scale interactions on
422	the seasonal modulation of ocean dynamics by the atmosphere, Nature communications,
423	5, 5636.
424	Shao, A. E., S. Mecking, L. Thompson, and R. E. Sonnerup (2013), Mixed layer satu-
425	rations of CFC-11, CFC-12, and SF ₆ in a global isopycnal model, J. Geophys. Res.

426 Oceans, 118, 4978–4988, doi:10.1002/jgrc.20370.

427	Simmons, A., S. Uppala, D. Dee, and S. Kobayashi (2007), ERA-Interim: New ECMWF
428	reanalysis products from 1989 onwards, in ECMWF Newsletter, vol. 110, pp. 25-35,
429	ECMWF.
430	Simmons, H. L., and M. H. Alford (2012), Simulating the Long-Range Swell of
431	Internal Waves Generated by Ocean Storms, Oceanography, 25(2), 30-41, doi:
432	http://dx.doi.org/10.5670/oceanog.2012.39.
433	Song, H., J. Marshall, P. Gaube, and D. J. M. Jr. (2015), Anomalous chlorofluorocarbon
434	uptake by mesoscale eddies in the Drake Passage region, J. Geophys. Res. Oceans, 120,
435	1065–1078, doi:10.1002/ 2014JC010292.
436	Song, H., J. Marshall, M. J. Follows, S. Dutkiewicz, and G. Forget (2016a), Source waters
437	for the highly productive Patagonian shelf in the southwestern Atlantic, J. Mar. Syst.,
438	158, 120–128.
439	Song, H., J. Marshall, D. R. Munro, S. Dutkiewicz, C. Sweeney, D. J. M. Jr., and
440	U. Hausmann (2016b), Mesoscale modulation of air-sea CO ₂ flux in Drake Passage,
441	J. Geophys. Res., 121, 6635-6649, doi:10.1002/2016JC011714.
442	Tagliabue, A., JB. Sallée, A. R. Bowie, M. Lévy, S. Swart, and P. W. Boyd (2014),
443	Surface-water iron supplies in the Southern Ocean sustained by deep winter mixing,
444	Nature Geoscience, 7, 314-320, doi:10.1038/ngeo2101.
445	Takahashi, T., J. Olafsson, J. G. Goddard, D. W. Chipman, and S. C. Sutherland (1993),
446	Seasonal variation of CO_2 and nutrients in the high-latitude surface oceans: A compara-
447	tive study, Global Biogeochemical Cycles, 7(4), 843-878.
448	Takahashi, T., S. C. Sutherland, R. Wanninkhof, C. Sweeney, R. A. Feely, D. W. Chipman,
449	B. Hales, G. Friederich, F. Chavez, C. Sabine, A. Watson, D. C. E. Bakker, U. Schus-
450	ter, N. Metzl, H. Yoshikawa-Inoue, M. Ishii, T. Midorikawa, Y. Nojiri, A. Körtzinger,
451	T. Steinhoff, M. Hoppema, J. Olafsson, T. S. Arnarson, B. Tilbrook, T. Johannessen,
452	A. Olsen, R. Bellerby, C. S. Wong, B. Delille, N. R. Bates, and H. J. W. de Baar
453	(2009), Climatological mean and decadal change in surface ocean pCO_2 , and net sea-
454	air CO ₂ flux over the global oceans, Deep-Sea Res. Pt. II, 56(8-10), 554-577.
455	Talley, L. D., G. Pickard, W. Emery, and J. Swift (2011), Descriptive Physical Oceanogra-
456	phy: An Introduction (Sixth Edition), 560 pp., Elsevier, Burlingham, MA.
457	Tian, F. (2016), Effects of coupling frequency on climate simulated by a coupled AO-
458	GCM, Ph.D. thesis, Universität Hamburg Hamburg.

- ⁴⁵⁹ Tulloch, R., R. Ferrari, O. Jahn, A. Klocker, J. LaCasce, J. Ledwell, J. Marshall, M.-J.
- 460 Messias, K. Speer, and A. Watson (2014), Direct estimates of lateral eddy diffusivity
- ⁴⁶¹ upstream of Drake Passage, J. Phys. Oceanogr., 44, 2593–2616.
- Verdy, A., S. Dutkiewicz, M. J. Follows, J. Marshall, and A. Czaja (2007), Carbon dioxide
 and oxygen fluxes in the Southern Ocean: Mechanisms of interannual variability, *Global Biogeochem. Cycles*, *21*, GB2020, doi:10.1029/2006GB002916.
- 465 Wanninkhof, R. (1992), Relationship between wind speed and gas exchange over the
- 466 ocean, J. Geophys. Res., 97(C5), 7373–7382, doi:10.1029/92JC00188.
- 467 Wetzel, P., A. Winguth, and E. Maier-Reimer (2005), Sea-to-air CO2 flux from
- ⁴⁶⁸ 1948 to 2003: A model study, *Global Biogeochem. Cycles*, 19, GB2005, doi:
- 469 10.1029/2004GB002339.