

GISS-E2.1: Configurations and Climatology

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Key Points:

- GISS E2.1 is an updated climate model version for use within the CMIP6 project.
- Atmospheric composition is calculated consistently in all model versions.
- Results demonstrate a very significant improvement in skill in a climate model without changes to atmospheric resolution.

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Abstract

This paper describes the GISS-E2.1 contribution to the Coupled Model Intercomparison Project, Phase 6 (CMIP6). This model version differs from the predecessor model (GISS-E2) chiefly due to parameterization improvements to the atmospheric and ocean model components, rather than increases in resolution. Model skill when compared to modern era climatologies is significantly higher than in previous versions. Additionally, updates in how forcings and forcing uncertainty are tackled have a material impact on the results. In particular, there have been specific improvements in representations of modes of variability (such as the Madden-Julian Oscillation and other modes in the Pacific) and significant improvements in the simulation of the climate of the Southern Oceans, including sea ice. The effective climate sensitivity to $2\times\text{CO}_2$ is slightly higher than previously at $2.7\text{--}3.1^\circ\text{C}$ (depending on version), and is a result of lower radiative forcing and stronger positive feedbacks.

Plain Language Summary

This paper describes and documents the latest iteration of the NASA GISS climate model which will be used for understanding historical climate change and to make projections for the future. We compare the model output to a wide range of observations over the recent era (1979–2014) and show that there has been a significant increase in how well the model performs compared to the previous version from 2014, though some persistent biases remain. The model has a temperature response to the increase of carbon dioxide that is slightly higher than it used to be, but is well within the range expected from observational and past climate constraints.

1 Introduction

The evaluation and assessment of climate models that are being used for attribution of past change and projections of future change has, for the last two decades, been dominated by the Coupled Model Intercomparison Project (CMIP). This is an internationally organised project run by the community and with almost universal participation from climate modeling groups across the world. The latest iteration (Phase 6) started accepting data in 2018 [Eyring *et al.*, 2016] in anticipation of the upcoming Intergovernmental Panel on Climate Change (IPCC) 6th Assessment Report (AR6) due in April 2021.

Climate modeling at the Goddard Institute for Space Studies (GISS) has a long pedigree dating back to the late 1970s [Hansen *et al.*, 1983, 1997, 2002] and has participated in almost all phases of the CMIP project, notably in CMIP3 and CMIP5 [Schmidt *et al.*, 2006, 2014]. Community experience over the last decade has demonstrated that constrained structural diversity in climate modeling is essential for elucidating important connections between processes and outcomes, and GISS models, with their distinct pedigree, have an important and continuing role to play in providing part of that diversity [Knutti *et al.*, 2013]. However, for that role to be successful, GISS needs to maintain and improve model realism (better process inclusion and higher skill) and continue participation in international and national climate model assessment projects. These projects allow model developers to benefit from the very broad scrutiny of results in these public archives from interested researchers and users across the world.

This paper is a description and an initial assessment of the GISS-E2.1 climate model, the first GISS contribution to CMIP6. This model version was developed as part of a long term strategy to improve model performance as much as possible without a significant jump in computational resources, building from the GISS-E2 models used in CMIP5. This exercise could be seen as the result of a much longer tuning process than is generally undertaken with a new model [Schmidt *et al.*, 2017]. Other GISS contributions to CMIP6 will include more significant leaps (higher vertical and horizontal resolutions, new grids and advection schemes, higher model top, new moist process schemes etc.) and will be reported in Rind *et al.* [2020] and elsewhere. Descriptions of specific results for the composition modeling,

80 historical runs and future projections will also be described elsewhere [*Bauer and Tsigaridis,*
81 2020; *Miller et al., 2020*].

82 1.1 Nomenclature

83 The series of GISS ModelE versions used in this and previous CMIP iterations, have
84 been GISS-E-R, GISS-E-H (in CMIP3, with the R and H denoting different ocean models
85 [*Schmidt et al., 2006; Hansen et al., 2007; Sun and Bleck, 2006*]) followed by GISS-E2-
86 R and GISS-E2-H in CMIP5 [*Schmidt et al., 2014*], and GISS-E2.1-G and GISS-E2.1-H
87 (in CMIP6). Other CMIP6 versions will be GISS-E2.2-G/H and GISS-E3-G. Some ver-
88 sions (denoted by -CC) also include an interactive carbon cycle [*Romanou et al., 2014*]. In
89 CMIP5, there were three formal versions of the models that varied according to the degree
90 of interactivity in atmospheric composition (`physics-version=1, 2, or 3`). In CMIP6,
91 `physics-version=2` has been dropped, `physics-version=1` denoted as NINT (for non-
92 interactive) uses offline ozone and aerosol fields from `physics-version=3` the OMA model
93 as described in [*Bauer and Tsigaridis, 2020*] and two new aerosol schemes have been added:
94 TOMAS (denoted by `physics-version=4`) [*Lee and Adams, 2012*] and MATRIX (`physics-version=5`)
95 [*Bauer et al., 2008*], which will be described elsewhere. Additionally, there is an additional
96 parameter in the CMIP6 database related to forcings. The `f#` number is used to denote dif-
97 ferent sets of forcing data (either concentrations, emissions or similar). In the initial E2.1
98 submission two versions have been made available for the historical runs; `f1` and `f2` which
99 have different ozone forcings (see section 2.1.3). Documentation of these conventions in all
100 GISS CMIP6 submissions will be maintained and updated at <https://data.giss.nasa.gov/modelE/cmip6/>.
101

102 2 Model code changes

103 Code changes since GISS-E2-R/H [*Schmidt et al., 2014*] consist of replacement or
104 structural variation of some parameterizations, updating of input files, bug fixes, and retun-
105 ing of specific parameters. These changes have been driven by internal and external identi-
106 fication of unsatisfactory performance, desired improvements in physical realism in param-
107 eterizations, and updates of observational data sets used either as input or evaluation. This
108 section lays out the drivers of specific changes and the specific changes made. Notably, with
109 the exception of additional layers in the ocean models (8 in E2.1-G, 6 in E2.1-H), no other
110 changes were made to the horizontal or vertical resolution in any component.

111 The main focus of the developments were to address unrealistic aspects in the CMIP5
112 simulations, notably poor Southern Ocean SST and sea ice (a common problem in across
113 CMIP5 [*Hyder et al., 2018*]), excessive ocean mixing, and precipitation pattern biases which
114 were evident in *Schmidt et al.* [2014]. Additionally, through the intense analysis by the wider
115 community of the CMIP5 simulations, additional issues were identified that led to subse-
116 quent bug fixes or re-calibrations of the code (for instance the assessment in *Prather et al.*
117 [2017] led to a reexamination of the ozone chemistry, and the authors of *Hezel et al.* [2012]
118 alerted us to an issue with snow cover over sea ice). Lastly, new functionality was required to
119 accommodate more complex emission input data and irrigation effects. The specifics of the
120 changes are outlined in the following sections.

121 2.1 Atmospheric processes

122 As stated above, atmospheric resolution is the same as in the CMIP5 model, including
123 the number of layers. However, a change was made to the manner in which terrain-following
124 (`sigma`) layers in the troposphere transition to constant-pressure layers in the stratosphere.
125 In E2, the transition is abrupt, occurring at 150 hPa. For E2.1, the option was activated to
126 use a smooth transition, centered at 100 hPa with a half-width of approximately 30 hPa. This

127 change removes some artifacts previously seen in the diagnostics but negatively impacts the
128 stratosphere circulation.

129 **2.1.1 Radiative Transfer**

130 The total solar irradiance has been updated based on new satellite calibrations [*Kopp*
131 *and Lean*, 2011] to have a base value of 1361 W m^{-2} (compared to 1366 W m^{-2} in GISS-
132 E2) though this is not expected to have any impact on the climatology or sensitivity once the
133 models have been retuned for radiative balance [*Rind et al.*, 2014]. Spectral irradiance values
134 have also been updated to the latest estimates [*Coddington et al.*, 2016].

135 Further calibration of the GISS-E2 radiation framework against line-by-line results led
136 to a few improvements for E2.1. Most notably, non-continuum absorption of shortwave radi-
137 ation by water vapor was significantly increased, thereby rectifying a problem subsequently
138 highlighted in analyses of the CMIP5 ensemble [*DeAngelis et al.*, 2015]. In the longwave re-
139 gion, a systematic increase of OLR of a few W m^{-2} was the main outcome of optimizations
140 of lookup tables for finer model layering and larger training sets of atmospheric profiles. HI-
141 TRAN 2012 spectroscopy [*Rothman et al.*, 2013] was also incorporated, with negligible im-
142 pact. The improvements to clear-sky SW and LW skill relative to E2 and other schemes can
143 be seen in the intercomparison of *Pincus et al.* [2015].

144 A small but consequential error in the snow masking of vegetation (where a constant
145 snow density was used instead of the computed predicted snow density) was fixed, thereby
146 reducing the area fraction of old, compacted snow and hastening springtime snowmelt.

147 **2.1.2 Clouds, convection and boundary layer**

148 As described in *Kim et al.* [2012], *Del Genio et al.* [2012] and *Del Genio et al.* [2015],
149 modifications to the cumulus parameterization in GISS-E2 led to much greater variability
150 associated with the Madden-Julian Oscillation (MJO) in GISS-E2.1. GISS-E2.1 retains the
151 basic entraining double plume updraft-downdraft framework used in GISS-E2, but with the
152 following changes: (1) The entrainment rate coefficient of the more weakly entraining plume
153 is increased from 0.3 to 0.4, thus increasing the sensitivity of convection to environmental
154 humidity; (2) The partitioning between convective precipitation that descends and has the
155 potential to evaporate in the environment rather than in the downdraft is increased from 0
156 percent to 50 percent, thus increasing the sensitivity of humidity to convection; (3) downdraft
157 buoyancy, which was determined solely by temperature in GISS-E2, is now based on virtual
158 temperature including condensate loading; (4) A previous limit on the cumulus mass flux
159 that inadvertently resulted in zero entrainment rates at high altitudes in strongly convecting
160 environments was eliminated.

161 The most impactful E2.1 update to the stratiform cloud parameterization concerns the
162 treatment of glaciation in the mixed-phase temperature range. In E2, glaciation in a given
163 gridcell is a probabilistically timed event after which no supercooled liquid can exist or form
164 until all ice has disappeared and the phase decision can "reset" for a new cloud. Within the
165 single-phase cloud condensate framework inherited from E2, E2.1 attempts to model glacia-
166 tion in a more continuous manner via a temperature-dependent autoconversion rate of super-
167 cooled liquid to ice precipitation. This rate is rapid at the homogeneous freezing temperature
168 of -35°C and decreases linearly toward the warm-cloud autoconversion rate at -5°C . Relative
169 to the new-cloud reset mechanism in E2, this "virtual" mixed-phase representation signifi-
170 cantly increases the amount of supercooled water cloud in the Southern Ocean and the Arctic
171 in E2.1. Radiation metrics, rather than liquid water path, were the primary means of evaluat-
172 ing this change.

173 The regime-specific threshold relative humidity for stratiform cloud formation in E2
174 was dependent upon moist convective activity, resolved vertical motion, and altitude (near
175 the surface). Convective area also restricted the maximum coverage of stratiform cloud. The

176 E2.1 code was modified as follows: (1) the coverage restriction is no longer applied above
 177 convective cloud top, (2) the dependence on vertical motion was dropped, since its applica-
 178 tion criterion did not distinguish fronts from other structures, and (3) altitude is taken to be
 179 relative to local planetary boundary layer (PBL) height rather than a fixed 850 hPa, better
 180 demarcating cloud-topped boundary layers from the free troposphere.

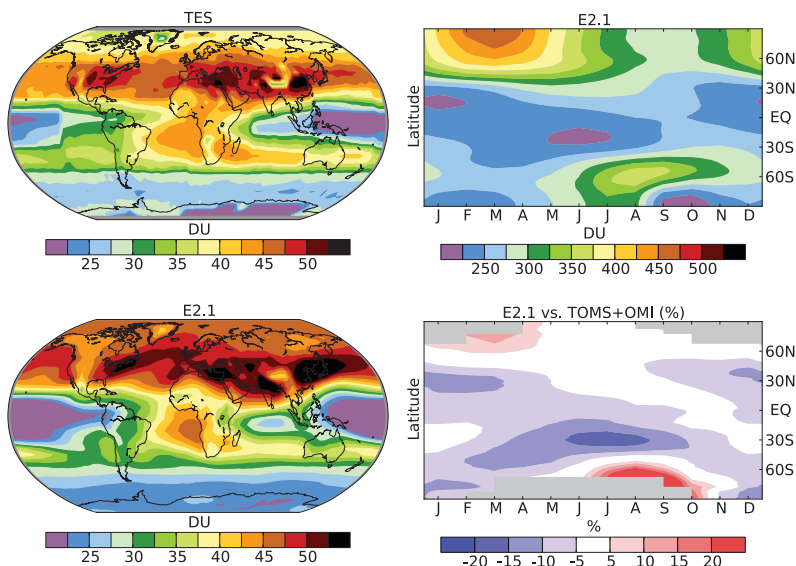
181 The modifications of the turbulence parameterization within and above the PBL [*Yao*
 182 *and Cheng*, 2012] from GISS-E2 include 1) the non-local vertical transport scheme for vir-
 183 tual potential temperature, specific humidity, and other scalars is updated from the [*Holt-*
 184 *slag and Moeng*, 1991] scheme to the more robust *Holtslag and Boville* [1993] scheme; 2)
 185 employing the turbulence length scale formulation obtained from the large eddy simulation
 186 data by *Nakanishi* [2001]; 3) using the more realistic "Richardson number criterion" rather
 187 than the "TKE criterion" to calculate the PBL height, following *Troen and Mahrt* [1986] and
 188 *Holtslag and Boville* [1993] and 4) modifying the similarity law near the surface in extreme
 189 stability conditions [*Zeng et al.*, 1998]. With the above modifications, the relative humidity
 190 and low cloud cover have better vertical structures due to greater transport of water vapor
 191 in the PBL. The differences in the diagnosed PBL height between the E2.1 and E2 versions
 192 correlate well with the differences in the total cloud distribution over oceans. This newer
 193 parameterization leads to improvement in cloud and radiation fields in the extratropics. Trop-
 194 ical low clouds were not specifically targeted, as they require finer layering at low levels and
 195 a cloud-enabled PBL scheme.

196 **2.1.3 Composition and chemistry**

197 The basic NINT simulations that are the focus of this paper do not have interactive
 198 composition, but the background fields of ozone and aerosol concentrations are derived from
 199 simulations of the interactive OMA version of the model, run under AMIP conditions [*Bauer*
 200 *and Tsigaridis*, 2020]. Thus the numerous, relatively minor updates and improvements to the
 201 composition modules affected these runs and so are described here for completeness.

202 All anthropogenic and biomass burning emissions of short-lived species were updated
 203 to CMIP6 specifications [*Hoesly et al.*, 2018; *van Marle et al.*, 2017], and are now prescribed
 204 annually, rather than by decadal interpolation as in CMIP5. Coding changes include: (1)
 205 calculating solar input to photolysis code using higher wavelength resolution; (2) updating
 206 the photolysis calculations to use up to 3 sets of temperature-dependent cross sections rather
 207 than 2; (3) harmonizing the heterogeneous chemistry reaction rate calculations in the strato-
 208 sphere to use the identical surface areas as those in the radiation code (typically satellite-
 209 derived extinction values); (4) correcting the passing of ozone from chemistry to radiation
 210 to use ozone only rather than all odd oxygen (which had led to warm biases in the upper-
 211 most stratosphere); (5) updating reaction rate coefficients from the JPL 2000 to the 2011
 212 compendium [*Sander et al.*, 2011]; (6) removing an imposed minimum tracer value which
 213 had led to large mixing ratios in high latitude grid boxes at high altitudes where total air
 214 masses are small; (7) expanding the representation of reactions including atomic hydrogen
 215 (no longer limited to specific pressure ranges); (8) expanding aircraft emissions to include
 216 more species; (9) correcting the amount of ozone input in photolysis calculations to use the
 217 gridbox top rather than the mid-gridbox value, which led to ozone chemistry biases [*Prather*
 218 *et al.*, 2017]. The harmonization of aerosol surface areas in (3) identified a coding error that
 219 led to large underestimates in volcanic aerosol surface areas for chemistry in the stratosphere.
 220 The two sets of runs denoted by f1 and f2 forcings reflect the impacts of that change.

221 Overall performance of the composition diagnostics is fairly similar to E2, based on
 222 comparison with the trace gas observations made in *Shindell et al.* [2013a]. A detailed anal-
 223 ysis suggests that over the US and China, the model is slightly high biased in terms of the
 224 simulated tropospheric ozone column relative to Tropospheric Emission Spectrometer (TES)
 225 measurements (Fig. 1) and substantially low biased in terms of aerosol optical depth relative
 226 to MISR observations [*Seltzer et al.*, 2017].



227 **Figure 1.** Left column: Annual average 2005–2009 tropospheric column ozone (DU) in TES observations
 228 (top) and in E2.1 (bottom). The tropopause is defined using the NCEP 2005–2009 monthly values for TES
 229 and the model’s internally calculated values for E2.1. Right column: 2000–2010 average of zonal mean,
 230 seasonal total column ozone (DU) in E2.1 (top) and the percent difference of this relative to TOMS + OMI
 231 observations for the same years ($100 \times (\text{model} - \text{obs}) / \text{obs}$) (bottom).

232 Several updates were made to lightning NO_x production in the chemistry module. The
 233 default flash rate parameterization remains a function of convective cloud depth, separately
 234 determined over land and sea [Price and Rind, 1994]. However, the calculation is now done
 235 using altitude above ground level rather than sea level, eliminating spurious lightning over
 236 high-altitude regions such as Antarctica. The land and marine flash rate equations are sepa-
 237 rately tuned to reproduce the respective present-day mean values from the Lightning Imaging
 238 Sensor (LIS) and Optical Transient Detector (OTD) satellite climatology [Cecil et al., 2014].
 239 Flash rates are converted to column NO_x production rates using a fixed NO_x -yield per flash
 240 assumption. These are then distributed vertically from the surface to the local cloud-top
 241 height using the unimodal probability distribution functions of Ott et al. [2010] instead of the
 242 earlier bimodal distribution of Pickering et al. [1998]. The NO_x -yield per flash is determined
 243 such the model reproduces the present-day methane chemical lifetime of 9.7 yr [Prather
 244 et al., 2012]. This results in 290 mol N per flash, yielding a global mean of 6.4 Tg N yr⁻¹.
 245 This is slightly lower than in E2 (7.3 Tg N yr⁻¹) [Shindell et al., 2013b] and falls within the
 246 relatively large range of estimates for the present-day lightning NO_x source (2–8 Tg N yr⁻¹)
 247 [Murray, 2016].

248 The aerosol module OMA [Bauer and Tsigaridis, 2020], simulating sulfate, nitrate,
 249 ammonium, carbonaceous aerosols (black carbon and organic carbon, including the NO_x -
 250 dependent formation of SOA and methanesulfonic acid formation), dust and sea-salt has re-
 251 mained mainly unchanged since CMIP5 (where it was called TCADI), with the exception of
 252 (1) increasing in-cloud ammonia dissolution to account for dissociation, thereby remedying
 253 the overabundance of nitrate aerosol in E2 [Nazarenko et al., 2017; Mezuman et al., 2016]
 254 (2) the parameterized e-folding time of hydrophobic to hydrophilic BC in OMA, a proxy
 255 for its aging lifetime, has been tuned in the CMIP6 OMA simulation to match that of MA-
 256 TRIX [Bauer et al., 2008], which does include physically-based aging calculations as part of

Pressure level (hPa)	Avg. diff. AMIP	Avg. diff. coupled	Avg. bias AMIP	Avg. bias coupled	Std. dev. of observations
125	43.5	65.8	9.9	45.1	92.9
200	21.7	27.3	1.2	7.7	52.2
300	13.4	15.2	7.0	8.3	25.6
500	9.6	10.9	6.2	7.5	11.7
900	8.0	8.8	3.6	4.5	8.9

258 **Table 1.** Ozone differences and biases (ppbv) between model E2.1-G versions and sonde climatologies.
259 Sonde data primarily from the 1990s and early 2000s; model from early 2000s.

257 the aerosol microphysics. The new aging timescale for OMA has been evaluated using ice
258 cores and HIPPO flight campaign data in [Bauer *et al.*, 2013]. And (3) updates to the dust
259 representation as discussed below. An evaluation of the aerosol scheme using CMIP6 forc-
260 ings against satellite, surface network and ice core data is discussed in [Bauer and Tsigaridis,
261 2020].

262 We updated the heterogeneous chemistry calculations for the formation of nitrate and
263 sulfate coatings on the surface of soil dust particles by uptake of nitric acid and sulfur diox-
264 ide, respectively, which were originally described by Bauer *et al.* [2004] and Bauer and Koch
265 [2005]. Dust properties are now retrieved from the dust module, instead of being defined
266 separately in the heterogeneous chemistry module, to make those properties consistent with
267 the rest of the model. This concerns the boundaries of the six dust bins (0.1–0.2, 0.2–0.5,
268 0.5–1, 1–2, 2–4, and 4–8 μm particle diameter), for which coatings on dust particles are cal-
269 culated in the model, the dust particle densities, and the weights that are used to partition
270 the dust mass of the total clay size range (0.1–2 μm), which is advected as a bulk in the model,
271 among the four clay bins. The weights reflect the size distribution of dust, compared to the
272 previous version where inadvertently only the largest clay bin was considered. An erroneous
273 calculation of the dust number concentration, which led to an overestimate of the number
274 was also corrected. The net effect of the changes are reduced masses of sulfate and nitrate
275 coating on dust by one order of magnitude due to lower uptake of the precursor gases sulfur
276 dioxide and nitric acid, respectively. The global masses of latter in the atmosphere are larger
277 by about 6% and 9%, respectively, with significantly larger increases over North Africa, Mid-
278 dle East, and Central Asia, where dust concentration is elevated. In turn, particulate nitrate
279 aerosols is up to five times higher over equatorial Africa and India and sulfate aerosol is up
280 to 50% higher in the northern hemisphere.

281 The default dust aerosol tracers in the OMA-version follow the approach of [Cakmur
282 *et al.*, 2006], with the difference that the emitted relative silt and clay fractions of total dust
283 and the emitted total dust mass are optimized in two successive steps, instead of simultane-
284 ously. The two-step approach reduces the emitted relative fraction of clay-dust mass (now
285 about 8% of all dust mass over the size range 0.1–32 μm for OMA), which makes the model
286 better conforming with recently published research on the global size distribution of dust in
287 the atmosphere [Kok *et al.*, 2017].

290 Ozone distributions used in the NINT models are generally similar to those in prior
291 versions. Changes to chemistry have resulted in modest improvements to comparisons with
292 observational data in the troposphere (Table 1). For example, the average bias near the sur-
293 face (900 hPa) has been reduced from 6.6 (22%) in E2 [Shindell *et al.*, 2013a] to 3.6 (12%)
294 in E2.1. This will help make for more realistic impact studies, which in the past required
295 substantial bias adjustments [Shindell *et al.*, 2018]. The modeled ozone hole in this config-
296 uration is biased as the Brewer-Dobson circulation to high-latitudes is too strong in winter,
297 leading to ozone and temperature overestimates during that season. This creates large pos-
298 itive biases in the lowermost stratosphere and upper troposphere from June through August

299 over 60-90S and smaller, but again positive, biases from December through April over 60-
300 90N.

301 Comparison of the tropospheric column ozone with observations from the Tropo-
302 spheric Emission Spectrometer (TES) show that the model captures many features of ozone's
303 distribution (Fig. 1). The effect of the wintertime positive biases in the lower stratosphere are
304 clearly visible in model overestimates of tropospheric columns poleward of 50°N and 70°S.
305 Such comparisons are highly sensitive to the tropopause definition [Shindell *et al.*, 2013a],
306 which is in turn sensitive to stratospheric temperature biases and so typically any widespread
307 ozone biases seen here reflect only small differences in the altitude of the tropopause rela-
308 tive to observations. The model captures the maximum over the Atlantic off the west coast
309 of Africa and the minima over the equatorial Pacific and Indian Oceans. As in E2, the min-
310 ima over the eastern tropical Pacific is too low, however, and this is likely to again dominate
311 biases in long wave radiative fluxes due to ozone [Bowman *et al.*, 2013]. The distribution of
312 column ozone is well represented over most of the NH mid-latitudes, though the magnitude
313 is roughly 2-4 DU too large. The global area-weighted average column in the model is 35.4
314 DU, very similar to the 35.9 DU from the TES observations. Spatial correlations are broadly
315 similar to those in E2, with an R^2 correlation against TES of 0.86 (compared to 0.85 in E2)
316 and a value of 0.68 against the tropospheric column estimate obtained from OMI minus MLS
317 observations (compared to 0.71 in E2).

318 **2.1.4 Gravity wave drag**

319 E2.1 includes orographic and frontal sources of parameterized gravity waves as in E2.
320 Systematic re-optimization of the scheme was not performed, but two updates required re-
321 calibration of tuning factors: (1) saturation momentum flux was reduced by a factor of ap-
322 proximately 2 as a result of correcting its definition (2) the metric for the presence of fronts
323 (deformation at 700 hPa) was corrected, increasing its magnitude. The orographic wave co-
324 efficient was thus reduced and the threshold deformation magnitude for generation of frontal
325 waves was increased.

326 **2.2 Ocean processes**

327 We used two ocean model versions with E2.1 which are denoted E2.1-G (coupling to
328 the GISS Ocean v1 (GO1)), and E2.1-H (coupling to HYCOM). This section describes the
329 updates in each since CMIP5.

330 **2.2.1 GISS Ocean v1**

331 For gross ocean structure and transport metrics, the most impactful updates to E2.1-G
332 are in the parameterizations of mesoscale eddies and vertical mixing. In addition, a high-
333 order advection scheme [Prather, 1986] and finer upper-ocean layering (an increase from
334 32 total layers to 40) sharpened the representation of frontal and thermocline structures in
335 regions of weak parameterized mixing. The updates outlined here will be described more
336 completely elsewhere, as part of parameter sensitivity studies.

337 A fundamental update to mesoscale eddy transport was the correction of an error in
338 the definition of neutral surfaces in E2-R which drastically reduced the restratification ef-
339 fect. Through the lens of ocean-only simulations and inter-model comparisons of temper-
340 ature/salinity drifts and circulation metrics such as AMOC and ACC strength, subsequent
341 work explored the consequences of controlled variations in the magnitude and structure of
342 the mesoscale eddy diffusivity [Marshall *et al.*, 2017; Romanou *et al.*, 2017]. Those efforts
343 informed the creation of a moderate-complexity 3D mesoscale diffusivity for E2.1-G whose
344 primary differences from the E2-R scheme are: (1) surface-intensified eddies, in the form
345 of an exponential decay of diffusivity with depth, where the location-dependent decay scale
346 is equal to $[\overline{|\rho_h z|}]/[\overline{|\rho_h|}]$, $[\overline{\cdot}]$ denotes vertical averaging, and ρ_h is the horizontal gradient of

347 potential density; (2) replacement of Rossby radius by a geographically constant nominal
 348 length scale $L = 39$ km in the $L^2 T_{\text{eddy}}^{-1}$ baroclinicity scaling of diffusivity retained from E2-R
 349 (J. Marshall, pers. comm.); (3) qualitative representation of the Coriolis element in the dis-
 350 carded Rossby radius by a factor $1/\max(.05, \sin(|\text{latitude}|))$ multiplying the diffusivity. The
 351 location dependence in (1) permits eddies to restratify the Southern Ocean over a large depth
 352 range, consistent with observed density structure there, while not overacting in other regions
 353 of the World Ocean (such as the North Atlantic, where the aforementioned ocean-only exper-
 354 iments indicated that deep mesoscale effects can suppress the AMOC). Simplifications (2)
 355 and (3) preserve the large-scale structure of the diffusivity distribution and its interactivity
 356 while eliminating unconstrained small-scale structure. E2.1-G also adopts a new representa-
 357 tion of mesoscale transport expressed in local quasi-isopycnal layering, circumventing some
 358 of the difficulties associated with the skew-flux representation that was employed in E2-R.

359 The E2.1-G vertical diffusivity now includes a contribution from tidal dissipation.
 360 AMOC sensitivity to this effect is exploited as a (model-specific) constraint on the consid-
 361 erable uncertainties surrounding this process. Exploratory coupled simulations, lacking the
 362 stabilizing effects of relaxation toward climatological surface salinity and a prescribed at-
 363 mospheric state, systematically developed a runaway haline stratification at high northern
 364 latitudes that was the proximate cause of a weak AMOC and excessive northern hemisphere
 365 sea ice. The sole parameterization change in any atmosphere or ocean component found able
 366 to sustain a strong AMOC was tidally driven mixing, which occurs in the shallow waters bor-
 367 dering the North Atlantic using the dissipation distribution generated by *Jayne* [2009].

368 Ventilation of marginal seas through their connecting straits has been increased via two
 369 mechanisms in E2.1-G, reducing salinity biases there. For straits deep enough that density
 370 contrasts can drive strong opposing flows at the surface and depth, the finer upper-ocean lay-
 371 ering in E2.1-G resolves this structure, in conjunction with a slight tuning of strait depths.
 372 Secondly, horizontal diffusivity was increased in straits that are shallow or have weaker den-
 373 sity contrasts. The first mechanism impacted the Red and Black seas, and the second the
 374 Baltic and Hudson. The first is the sole ventilation mechanism for straits narrower than the
 375 nominal resolution, which are parameterized using the *Russell et al.* [1995] 1-dimensional
 376 channel scheme that lacks horizontal mixing.

377 2.2.2 HYCOM

378 HYCOM is a hybrid-isopycnal ocean model that was used with previous coupled Mod-
 379 elE versions [*Sun and Bleck, 2006; Romanou et al., 2013*]. Compared to the version used in
 380 E2, for E2.1-H, the number of vertical layers was increased to 32 from 26, and the refined
 381 equatorial mesh was removed (since it no longer provided a demonstrable increase in skill).
 382 HYCOM has traditionally used σ_2 as its vertical coordinate: potential density referenced to a
 383 pressure nominally corresponding to 2km depth. At pressures far from this reference, stable
 384 in-situ stratification may be misdiagnosed as unstable according to potential density, impact-
 385 ing the layering scheme and vertical mixing. To ensure a monotonic potential density profile
 386 in the upper ocean under conditions of stable in-situ stratification there, E2.1-H employs σ_1
 387 (potential density referenced to 1km). This change eliminated spurious deep convection in
 388 the Southern Ocean which inhibited formation of the summer halocline and limited sea ice
 389 extent. The resulting degradation of the abyssal diagnosis of stratification was found to be
 390 benign.

391 The virtual salt flux formulation of surface freshwater fluxes, employed by HYCOM
 392 for consistency with its barotropic/baroclinic mode-splitting scheme, was corrected to con-
 393 serve global salt, thereby eliminating a net source that resulted in significant positive biases
 394 in E2-H salinity. Other fixes to ocean-atmosphere-ice flux coupling include (1) interpolation
 395 between grids, (2) elimination of slight inaccuracies in the sea ice mass and heat fluxes, and
 396 (3) a modification to the land topography along the coastline to reduce flux biases in atmo-
 397 spheric grid-boxes with average land heights significantly above sea level.

398 2.3 The cryosphere

399 Common to both ocean models as in E2, the sea ice component of E2.1 retains the
 400 overall framework of E2, excepting the treatment of salt as a material constituent. Algorithmic
 401 changes within the framework made the most direct contributions to differences with E2
 402 climatology, and include (1) correction of an inadvertent snow-to-ice transformation during
 403 vertical regridding, thereby increasing snow thickness and surface albedo, (2) removal of a
 404 10% floor on lead fraction for conditions typical of the Antarctic winter, (3) closure of leads
 405 for thick-ice conditions typical in the Arctic, thereby reducing wintertime heat flux and ice
 406 growth there, (4) independent horizontal advection of snow mass. Thermodynamics now
 407 follows the "Brine Pocket" (BP) parameterization [Bitz and Lipscomb, 1999; Schmidt et al.,
 408 2004], and thus salt plays a more active role in E2.1 sea ice, affecting its specific heat and
 409 melt rates. Processes relevant to the salt budget (e.g. gravity drainage and flushing of melt-
 410 water) are consistently processed with the BP physics. The switch from the previous 'Saline
 411 Ice' thermodynamics in E2 to the BP one in E2.1 lead to a slight increase in multiyear sea ice
 412 thickness and of sea ice area in the Arctic, a slight reduction of the Antarctic sea ice area as
 413 well as a more physically realistic vertical profile of the salinity in the ice.

414 2.4 Land surface processes

415 2.4.1 Irrigation and Groundwater

416 While transient historical changes to irrigation was implemented as a forcing in E2
 417 [Cook et al., 2014; Krakauer et al., 2016], it was not included in the standard CMIP5 sub-
 418 missions. In E2.1, irrigation is now a standard component. Seasonal irrigation amounts from
 419 Wisser et al. [2010] are applied to the grid box land surface. The water is drawn first from
 420 the local surface water system (including rivers and lakes), and if that is insufficient, it is as-
 421 sumed to be drawn from an external groundwater source (which is tracked diagnostically).
 422 Groundwater is assumed to have the same temperature as the soil, and has default tracer val-
 423 ues. Groundwater recharge is not accounted for, and so there is a small increase in total water
 424 mass (and eventually, sea level) associated with the net global groundwater draw in these
 425 simulations. These affects have a complex impact on freshwater delivery to the oceans (and
 426 hence sea level). Irrigation from local surface water sources, leads to increased soil moisture
 427 and reduced river outflow, while net removals of groundwater add freshwater to the climate
 428 system, to the tune of about 0.2 mm yr^{-1} of global sea level equivalent in 2010.

429 2.4.2 Vegetation

430 As in E2, all vegetation properties affecting physical climate, with the exception of
 431 canopy conductance, are prescribed in the simulations described here, whose primary up-
 432 date was the incorporation of satellite-derived distributions of vegetation characteristics, as
 433 described below. Like E2, E2.1 sees vegetation properties via the Ent Terrestrial Biosphere
 434 Model (Ent TBM), a demographic dynamic global vegetation model (DGVM) whose func-
 435 tionalities are gradually being coupled to ModelE [Kiang, 2012; Kim et al., 2015; Kiang
 436 et al., in prep.], including carbon cycle interactivity (Ito et al., in prep.). Prescribed interan-
 437 nual variation of vegetation is limited to land use and land cover change (LULCC), by which
 438 historical crop and pasture cover is used to rescale the natural vegetation cover fractions in a
 439 grid cell. CMIP6 protocol implementations are described more fully in Miller et al. [2020]
 440 and Ito et al. (in prep.).

441 The Ent Global Vegetation Structure Data set (Ent GVSD) [Kiang et al., in prep.] has
 442 been created as part of ongoing Ent TBM development and is now used for E2.1, replac-
 443 ing E2 prescriptions of vegetation cover, type, height, and LAI based on Matthews [1983].
 444 Ent GVSD satellite data sources include land cover types and monthly varying LAI from the
 445 Moderate Resolution Imaging Spectroradiometer (MODIS) [Gao et al., 2008; Myneni et al.,
 446 2002; Tian et al., 2002a,b; Yang et al., 2006], and tree heights from Simard et al. [2011],

447 who utilized 2005 data from the Geoscience Laser Altimeter System (GLAS) aboard the
 448 ICESat (Ice, Cloud, and land Elevation Satellite). Specific leaf area (mass carbon per leaf
 449 area) data from the TRY database of leaf traits [Kattge *et al.*, 2011] was classified for the
 450 Ent TBM 13 plant functional types (PFTs). These observed spatial distributions and leaf
 451 trait parameters together allow an estimate of vegetation demography (density of plants per
 452 area), and achieve equilibrium behavior in plant-atmosphere carbon exchange and internal
 453 plant carbon balances for late 20th C. to early 21st C. climate. The water stress algorithm,
 454 which controls the availability of soil water for transpiration, was replaced in E2.1 with a
 455 more commonly-used soil water deficit-based one [Porporato *et al.*, 2001; Rodriguez-Iturbe,
 456 2000], with the goal of improving transpiration, by distinguishing soil moisture levels at
 457 which onset of water stress happens for different plant functional types.

458 The overall effect of these updates upon surface albedo was significant in some re-
 459 gions, though the overall impact upon physical climate modest compared to other compo-
 460 nents. Ent PFTs are mapped to the E2 vegetation types for radiative purposes in E2.1; reclas-
 461 sification of cover types directly increased the surface albedo of Australia and eastern South
 462 America by several percent. High northern latitudes became brighter via increased snow
 463 masking, though this effect was compensated by the correction to snow masking described in
 464 the radiation section. Canopy conductances generally decreased using the new LAIs.

465 **3 Simulation design and configurations**

466 The GISS models are designed so that any experiment can be run with an appropri-
 467 ate level of interactivity and complexity - some experiments require the aerosol and chem-
 468 istry fields to respond to and influence the surface climate, while other simulations focus
 469 on one-way impacts. In earlier iterations, NINT historical simulations relied on calculated
 470 concentrations of aerosols and tropospheric ozone from a prior generation of models. For in-
 471 stance, the NINT simulations in CMIP5 (using GISS-E2-R or GISS-E2-H) used fields from
 472 Koch *et al.* [2011] which were calculated using the CMIP3 model (GISS-E). In CMIP3, the
 473 aerosol and ozone fields were from the SI2000 version of the model [Koch, 2001; Koch *et al.*,
 474 1999] and thus were not strictly consistent with the composition changes generated in the
 475 same-generation interactive models (OMA or MATRIX aerosol microphysical versions) or
 476 the specified emission paths. Additionally, many key interactions present in the (computa-
 477 tionally expensive) interactive runs (such as ozone responses to volcanoes or solar activity
 478 changes) were not represented in the CMIP5 NINT runs.

479 For CMIP6 we have striven for an increased coherence between forcings and model
 480 physics. Namely, we have generated all the historical composition fields using an ensemble
 481 of AMIP runs (1860–2014) with the updated interactive OMA version and annually-resolved
 482 CMIP6 emissions [Bauer and Tsigaridis, 2020]. The time needed to generate new compo-
 483 sition fields slows down production, but the resulting NINT simulations have more fidelity
 484 to the real world and reflect more processes, while being 3–4 times faster to run when com-
 485 pared to interactive composition versions.

486 **3.1 Pre-industrial boundary conditions**

487 There are a few notable changes from CMIP5 for “pre-industrial” (PI) conditions,
 488 which is a slight misnomer, since conditions around 1850 cannot be considered to be unaf-
 489 fected by industrialization, agriculture and fossil fuel use (through the background green-
 490 house gas levels) and explicit background levels of land use and land cover change, including
 491 irrigation [Hawkins *et al.*, 2017]. We now include a background level of irrigation along with
 492 background levels of land use/land cover (LULC) alterations and anthropogenic aerosols.
 493 The emissions from biomass burning are taken from the standard CMIP6 specifications, but
 494 include an (uncertain) anthropogenic component. The spin-up under PI conditions is always
 495 greater than 500 years, but the procedure does not include pre-1850 transient changes that
 496 might be expected to still have been responsible for ocean heat content anomalies at that time

497 [Stenchikov *et al.*, 2009; Gregory, 2010]. Nonetheless, the difference in sub-surface ocean
 498 conditions from reality in 1850 are significantly larger than the impact of prior transient vol-
 499 canic effects (compared to a suitable averaged background level). Experience from simula-
 500 tions of the last millennium in CMIP5 suggests that the differences in 20th Century transient
 501 climate resulting from this choice are minimal.

502 3.2 Historical Transients

503 As mentioned above, radiatively active atmospheric composition (ozone and aerosols)
 504 is taken from AMIP experiments using E2.1 (OMA) using CMIP6-prescribed annual emis-
 505 sions of aerosols, their precursors and other short-lived reactive chemical species. Well-
 506 mixed greenhouse gases, solar activity changes (affecting TSI and the spectral irradiance),
 507 LULC (including irrigation), were specified using a mix of approaches [Miller *et al.*, 2020].
 508 Volcanic aerosols were prescribed using pre-computed aerosol depth and effective particle
 509 radius [Thomason *et al.*, 2018], though we will also be using interactive emission-driven vol-
 510 canic effects in some future CMIP6 simulations [LeGrande *et al.*, 2016].

511 It is important to note that there is substantial uncertainty in some of these drivers over
 512 time, especially in the aerosols, solar activity, and early big volcanic eruptions. We therefore
 513 plan to incorporate this uncertainty in the CMIP6 historical simulations. Different realiza-
 514 tions of the forcings suite are denoted by the f number in the r1pf designation of each indi-
 515 vidual run in the CMIP6 archive. The first submissions with E2.1-G/E2.1-H have two varia-
 516 tions ($f1$ and $f2$) which differ due to a correction of the stratospheric ozone chemistry code
 517 in the OMA simulations (see section 2.1.3). The $f2$ simulations have an improved strato-
 518 spheric ozone response to volcanoes and a better representation of the Antarctic ozone hole,
 519 but are not detectably different in tropospheric climate or response.

520 4 Coupled Model Tuning

521 Model tuning for E2.1 followed the procedure described in Schmidt *et al.* [2017]. In
 522 the atmosphere-only simulation using the $f1$ 1850 pre-industrial conditions, the parameters
 523 in the cloud schemes that control the threshold relative humidity and the critical ice mass for
 524 condensate conversion are used to achieve global radiative balance and a reasonable global
 525 mean albedo. Tuning of convective entrainment was also used to enhance MJO variability
 526 [Del Genio *et al.*, 2015].

527 Upon coupling the ocean and atmosphere models, there is an initial drift to a quasi-
 528 stable equilibrium which is judged on overall terms for realism, including the overall skill in
 529 the climatological metrics for zonal mean temperature, surface temperatures, sea level pres-
 530 sure, short and long wave radiation fluxes, precipitation, lower stratospheric water vapor, and
 531 seasonal sea ice extent. For the configuration to be acceptable, drifts have to be relatively
 532 small and quasi-stable behavior of the North Atlantic meridional circulation and other ocean
 533 metrics, including the Antarctic Circumpolar Current, are required. ENSO-related metrics
 534 are also monitored, but they were not specifically tuned for, since the underlying tropical Pa-
 535 cific SST climatology was not considered to be a feasible tuning target using E2.1 vertical
 536 resolution, cloud, and boundary layer schemes. In practice, longer spin-up integrations help
 537 reduce drift, and the model state once stabilized can be assessed for suitability. Large drifts
 538 at the start of an integration have often been reduced by different tuning choices that either
 539 affect surface atmospheric fluxes or (more usually) ocean mixing.

540 Composition tuning (as described below) is done with pre-industrial and present-day
 541 fully interactive simulations (including chemistry and aerosols and indirect effects) and the
 542 non-interactive versions use the composition derived from those simulations and the same
 543 tuning. Specifically, in the NINT transient simulations, the AIE was preset to have a value
 544 of -1 Wm^{-2} in 2000 as it was in the CMIP5 simulations [Miller *et al.*, 2014]. Configurations
 545 with interactive aerosols have free latitude to produce whatever forcing is calculated. Ad-

ditionally, in previous model versions we used a temperature threshold for the formation of polar stratospheric clouds (and hence the heterogeneous chemistry associated with them) [Shindell *et al.*, 2013a] which was tuned so that the polar ozone hole timing is correct despite potential biases in polar vortex temperatures. However, this was not used in E2.1. This model does, however, maintain prior practice of tuning photolysis rates at short wavelengths (<200 nm) for N₂O and O₂ that corrects for problems in stratospheric circulation that otherwise lead to biases in high latitude concentrations of NO_x and O₃. Also, with respect to dust aerosols, emissions are tuned so that the model can match retrieved aerosol optical depths for the present-day [Miller *et al.*, 2006], similarly tuning of the lightning parameterization (and associated source for NO_x) is done against modern observations of flash rate, and tropospheric ozone amounts.

Note that the atmospheric component was tuned using the pre-industrial f1 background ozone and aerosols. Upon switching to the f2 background, there was a slight drift in the coupled model. Prior to any historical runs with the f2 forcings, the coupled model was run a further 100 years to reach a new quasi-equilibrium.

We do not fine tune for an exact global mean surface temperature, since that is effectively precluded by the long spin-up times and limited resources available. Similarly, no tuning was done for climate sensitivity or for performance in a simulation with transient forcing or hindcasts.

5 Climatology 1979–2014

As was seen in the results shown in Schmidt *et al.* [2014], the impact of interactivity in the aerosol or chemistry parts of the model have limited impacts on the climatologies. In addition, while in E2, there was a substantive difference in the composition fields between NINT and TCADI simulations, that is no longer the case in E2.1, though composition-related interactivity may have a greater impact on the variability. We therefore only show the ensemble mean climatology from the standard NINT simulations (10 members for E2.1-G, 5 members for E2.1-H), in both spatial patterns, zonal and global means compared to updated observed climatologies for the satellite period (or as close as possible). All diagnostics are from the f2 historical simulations unless otherwise stated. We include the zonal mean diagnostics from the E2.1-G f1 forcings ensemble for completeness, but the differences are small. Note that the map projection uses Equal Earth [Šavrič *et al.*, 2018] and that we now plot zonal means with an area weighted x-axis to minimise visual distortion.

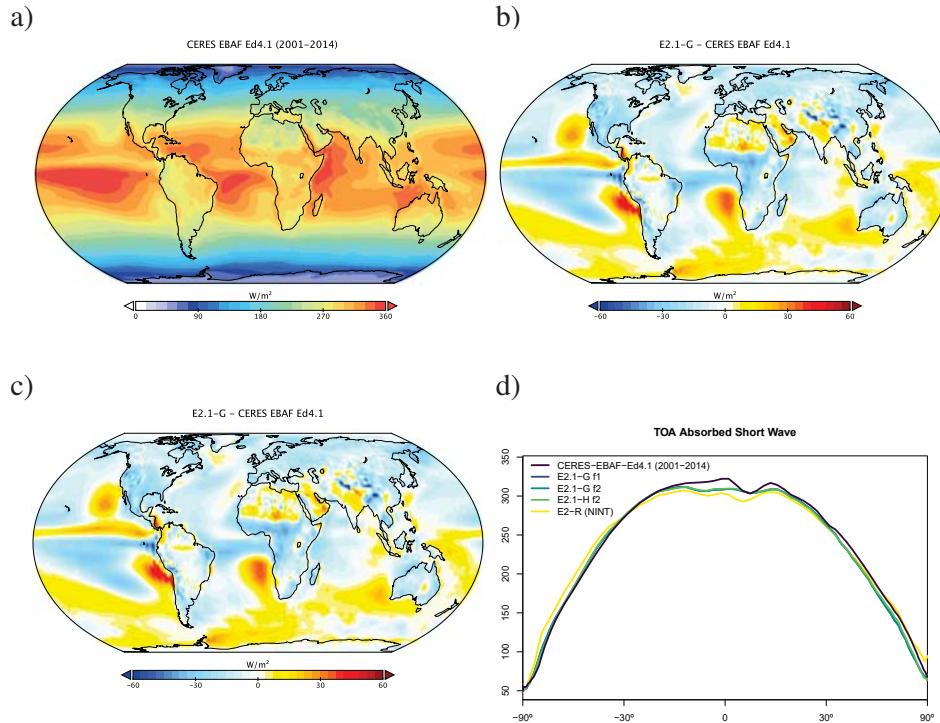
5.1 Global mean diagnostics

Table 2 summarizes a standard set of global mean diagnostics for the NINT versions of the GISS-E2.1 models (with f2 forcings) and a comparison with updated observations and previous model versions [Schmidt *et al.*, 2014]. Notable improvements are in the global mean temperature, precipitation, and sensible heat fluxes. The net radiative imbalance over this period is also in better comparison with updated estimates from NODC. There are notable biases in total column water vapour (7% too high), and LW cloud forcing (some 20 to 25% too low, though still better than previously). Lower stratospheric water vapour is deficient, consistent with a too cold tropopause. The TOA radiative fluxes are tuned for in pre-industrial atmosphere-only simulations and are therefore not truly predictive. Differences between the coupled models with different ocean modules are small compared to differences with the observations at the global mean level.

Field	E2.1-G	E2.1-H	E2-R	E2-H	Observations
Surface air temp. (°C)	14.1	14.5	14.9	15.6	14.3±0.5 ^J
Planetary Albedo	30.4	30.2	29.9	29.7	30 ^E /29.4 ^{SEA}
Cloud cover (%)	59.9	59.8	62	62	68 ^{SRK}
Precip. (mm day ⁻¹)	2.97	2.98	3.17	3.21	2.9 ^G
Snowfall (mm day ⁻¹)	0.24	0.23	0.19	0.17	0.18 ^{L08} /0.12 ^{SEA}
Atmos. water (mm)	26.7	26.8	23.8	24.0	24.9 ^O
Energy fluxes (W m ⁻²):					
TOA Absorbed SW	236.9	237.5	239.5	240.3	240.2 ^{SEA} /239.4 ^T
TOA Outgoing LW	236.5	237.1	238.8	239.5	239.7 ^{SEA} /238.5 ^T
Surf. Abs. SW	161.5	161.9	169.5	170.1	165 ^{SEA} /169 ^T
Surf. Down. LW	345.8	347.4	341	344	345.6 ^{SEA} /343 ^T
Surf. Net LW (up)	50.5	50.7	56.9	56.9	52.4 ^{SEA} /57 ^T
Sensible heat flux	23.9	23.9	19.3	19.0	24 ^{SEA} /17 ^T
Latent heat flux	85.8	86.2	91.9	92.8	88 ^{SEA} /82 ^T
TOA SW cld. forcing	-48.8	-48.1	-48.9	-48.5	-45.4 ^C
TOA LW cld. forcing	21.1	21.1	18.8	19.0	25.9 ^C
TOA Net. Rad. Imb.	0.42	0.39	0.66	0.62	0.41±0.03 ^{NN}
Trop. lower strat. water vapor minima (ppmv)					
	3.0	2.8	4.5	4.4	3.8±0.3 ^D
Zonal mean tropopause temp. (min., DJF) (°C)					
	-81	-82	-80	-80	-80
Hadley Circ. (10 ⁹ kg s ⁻¹) (DJF)					
	205	207	206	208	170–238 ^S

579 **Table 2.** Global annual ensemble mean model features over the period 1979-2014 (1980-2004 for the E2
580 models) and key diagnostics compared to observations or best estimates. Cloud cover is estimated based on
581 clouds with optical thickness >0.1 . ^J Jones *et al.* [1999] with updates, ^C CERES EBAF Ed4.1 Loeb *et al.*
582 [2019], ^T Trenberth *et al.* [2009] and updates, ^G GPCP V2.3/TRMM TMPA V7 Huffman *et al.* [2007, 2009],
583 ^O Obs4MIPs, ^{NN} Dervied from NOAA NCEI ocean heat content data, ^D Dessler [1998], ^S Stachnik and
584 Schumacher [2011], ^{SEA} Stephens *et al.* [2012], ^{SRK} Stubenrauch *et al.* [2013]

TOA Absorbed Solar Radiation



597 **Figure 2.** a) Annual climatology of TOA Absorbed Short Wave (W m^{-2}) in CERES EBAF Ed4.1 [Loeb
 598 *et al.*, 2019]. b) Difference of E2.1-G from CERES. c) Difference from E2.1-H. d) Absolute Zonal means,
 599 including E2.1-G (f1 and f2), E2.1-H and the earlier model version, E2-R.

5.2 Radiation and Clouds

596

601 Radiation diagnostics are now compared to the latest balanced CERES product (EBAF
 602 Ed4.1) Loeb *et al.* [2019]. Improvements since E2 are clearest in the Southern Ocean, where
 603 excessive SW absorption has been greatly ameliorated, but also in the tropics, although ob-
 604 vious biases associated with the marine stratus regions in the eastern ocean basins still exist
 605 (figs. 2 and 3). Notably the sign of the biases in the Arctic have changed in SW absorption.
 606 There is a lack of asymmetry across the tropics (which is clear in the observations), with the
 607 southern tropics characterised by excessive water vapor and cloud forcing. Note that South-
 608 ern Ocean estimates of TOA absorbed SW (fig. 2) are somewhat better constrained than SW
 609 cloud radiative forcing (fig. 7).

610 Cloud fraction diagnostics have been upgraded to the ISCCP-H product for 1984–
 611 2014 [Young *et al.*, 2018]. The overall patterns are slightly improved in the tropics and mid-
 612 latitudes, but the persistent biases (in the marine stratus regions) remain clear (figs. 4 and 5).
 613 The bias in low cloud over sea ice regions may be an artifact. The improvements are clearer
 614 in the SW CRF diagnostic (fig. 7), and in the high latitudes at least for the LW cloud radiative
 615 forcing which remains overall too low (except in the erroneously cloudy tropical mid-Pacific
 616 (fig. 8). The cloud top pressure/cloud optical depth histograms (fig. 6) show that the model
 617 has improved its “too few - too bright” low cloud problem, as low cloud cover has increased
 618 and optical thickness has decreased in relation to the E2 version.

TOA Outgoing Long Wave Radiation

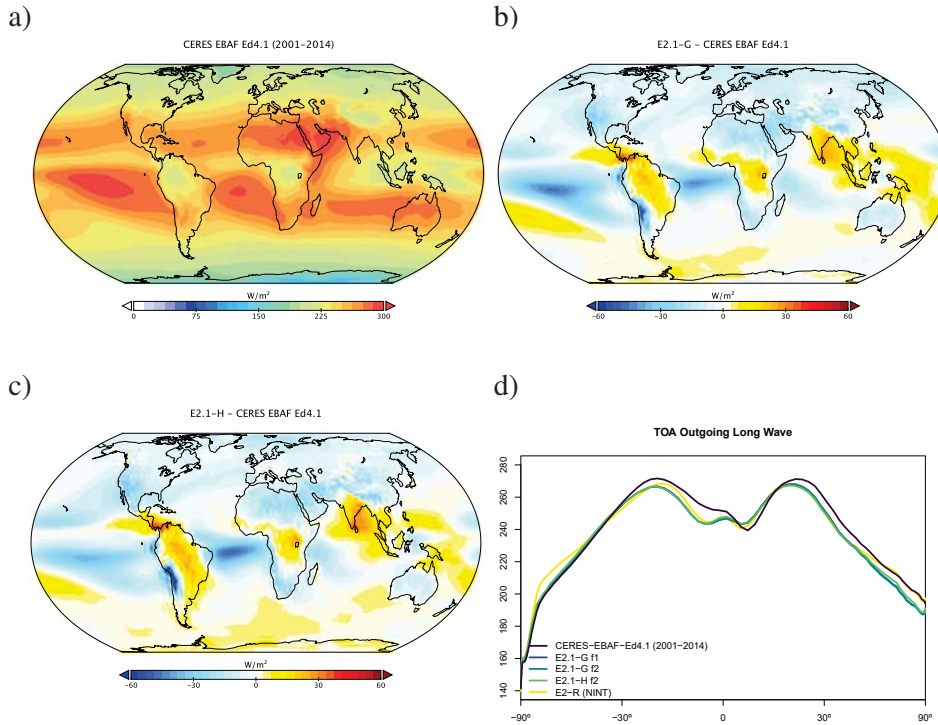
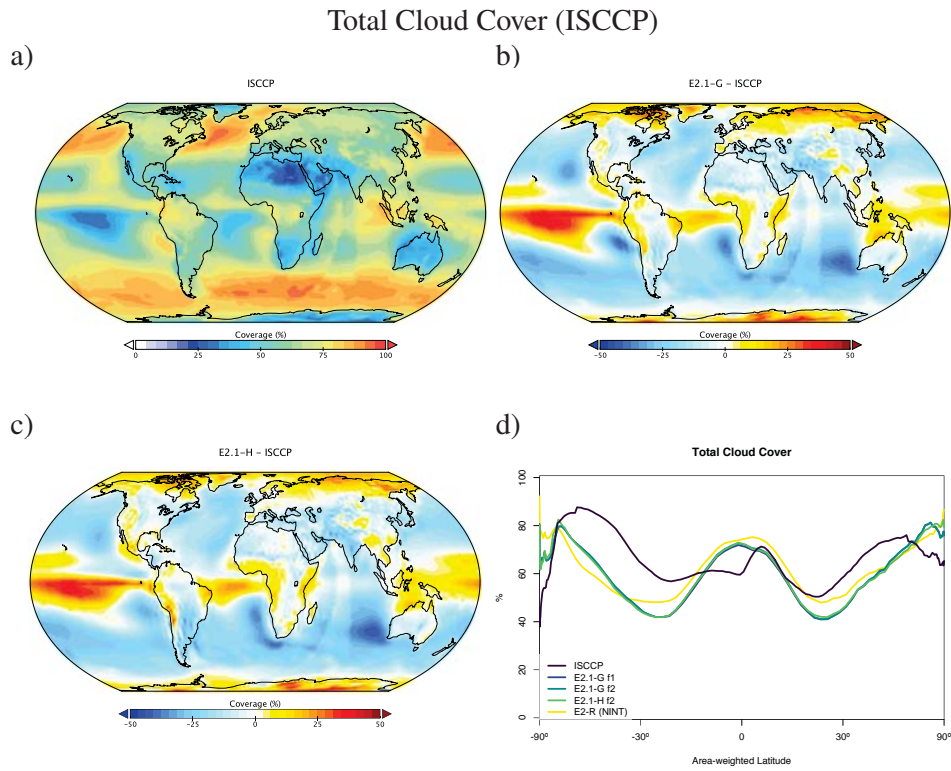


Figure 3. Annual climatology of TOA outgoing long wave data in data and models, as in fig. 2.

600

619 Additional information from comparisons of an earlier E2.1 version with active-sensor
 620 satellite observations (not shown) confirms an improvement of the low cloud cover in the
 621 high latitudes and over the trade wind regions while large biases remain over the stratocumulus
 622 regions in the tropics and subtropics. This low cloud bias might alter the strength of the
 623 low cloud feedbacks in response to global warming [Cesana *et al.*, 2019]. The large high-
 624 cloud positive bias found in E2 [Cesana and Waliser, 2016] has been mostly removed except
 625 in the southern tropics, where the overestimate of total cloud cover (fig. 4) comes from an
 626 excess of very high clouds (above 16 km), which are not present in satellite observations. Fi-
 627 nally, the amount of supercooled water cloud with respect to ice cloud is overestimated on
 628 average [Cesana *et al.*, 2016]. As a result, the negative cloud-phase feedback [Ceppi *et al.*,
 629 2016; Tan *et al.*, 2016] might be underestimated in E2.1 while it was overestimated in E2.

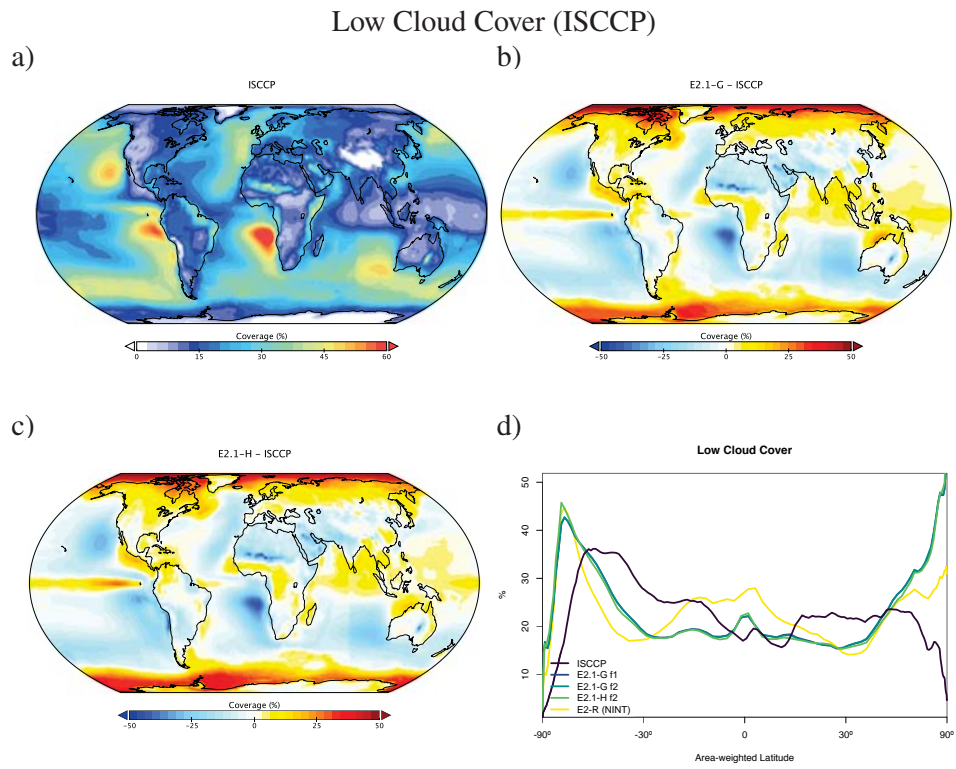
630 Atmospheric hydrological observations come from two blended data products via
 631 the Obs4MIPS archive [Gleckler *et al.*, 2011; Teixeira *et al.*, 2014; Ferraro *et al.*, 2015].
 632 The precipitable water vapor is a blend of the RSS product over ocean [Wentz and Schabel,
 633 2000; Wentz *et al.*, 2007] and MERRA-2 (over land) from the CREATE-MRE project [Pot-
 634 ter *et al.*, 2018] while the precipitation product is a blend of TRMM satellite estimates over
 635 ocean [Huffman *et al.*, 2007; Adler *et al.*, 2009] and GPCP [Huffman *et al.*, 2009] Version
 636 2.3 satellite-gauge calibrated precipitation over land. Precipitable water vapor discrepan-
 637 cies (fig. 9) are larger than in the previous model in the tropics, where the lack of asymmetry
 638 is readily apparent. The largest biases in water vapor coincide with the excessive LW CRF.
 639 This is also consistent with overall precipitation biases (fig. 10) which show a classic double-
 640 ITCZ problem in the Pacific, although one that is diminished in magnitude compared to E2.
 641 Excessive land precipitation in the Western Pacific Warm Pool has also been greatly ame-



646 **Figure 4.** Annual climatology of Total Cloud Cover as seen by ISCCP-H, figure description as in fig. 2.

642 liorated. Note too, that part of the reduced bias in rainfall is due to upgrades in the observa-
 643 tional product.

644 Snowfall biases are noticeable in the zonal mean (fig. 11), particularly in the Arctic,
 645 where excessive snowfall is related to wintertime cold biases in both models.

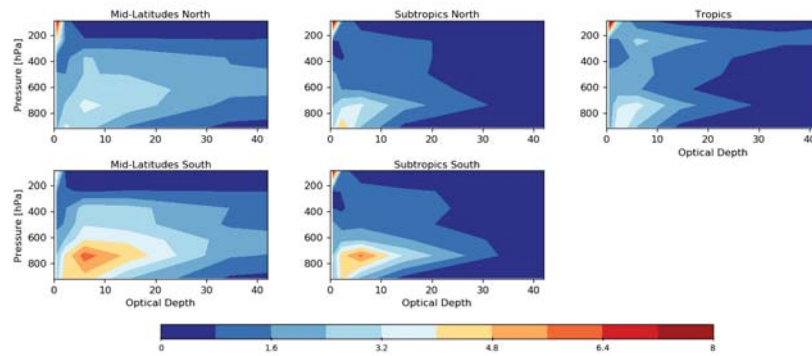


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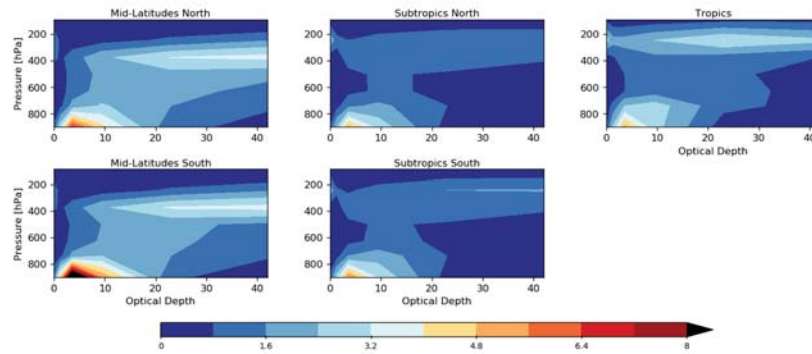
Figure 5. Annual climatology of Low Cloud Cover as seen by ISCCP-H, figure description as in fig. 2.

Regional Cloud Climatology

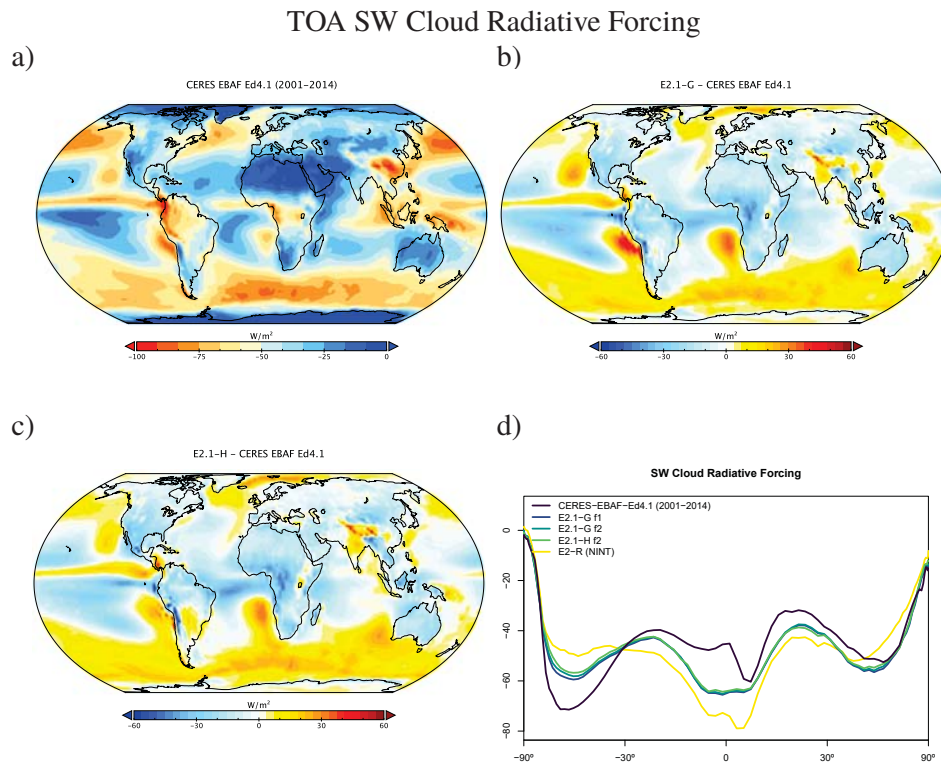
a) ISCCP-H



b) E2.1-G

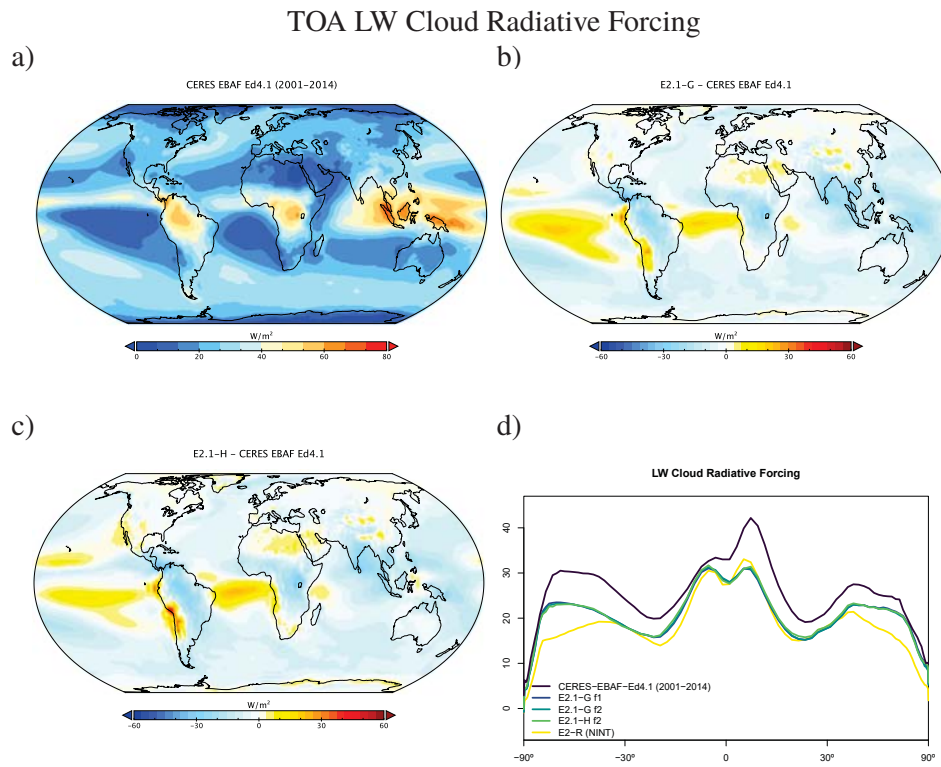


648 **Figure 6.** Climatology of cloud occurrence as a function of optical depth and pressure for five latitudinal
 649 bands as seen by ISCCP (60°N–30°N, 30°N–15°N, 15°N–15°S, , 15°S–30°S and 30°S–60°S). a) Data from
 650 ISCCP-H [Young *et al.*, 2018]. b) Data from the ensemble mean E2.1-G results. (Results from E2.1-H are
 651 indistinguishable).



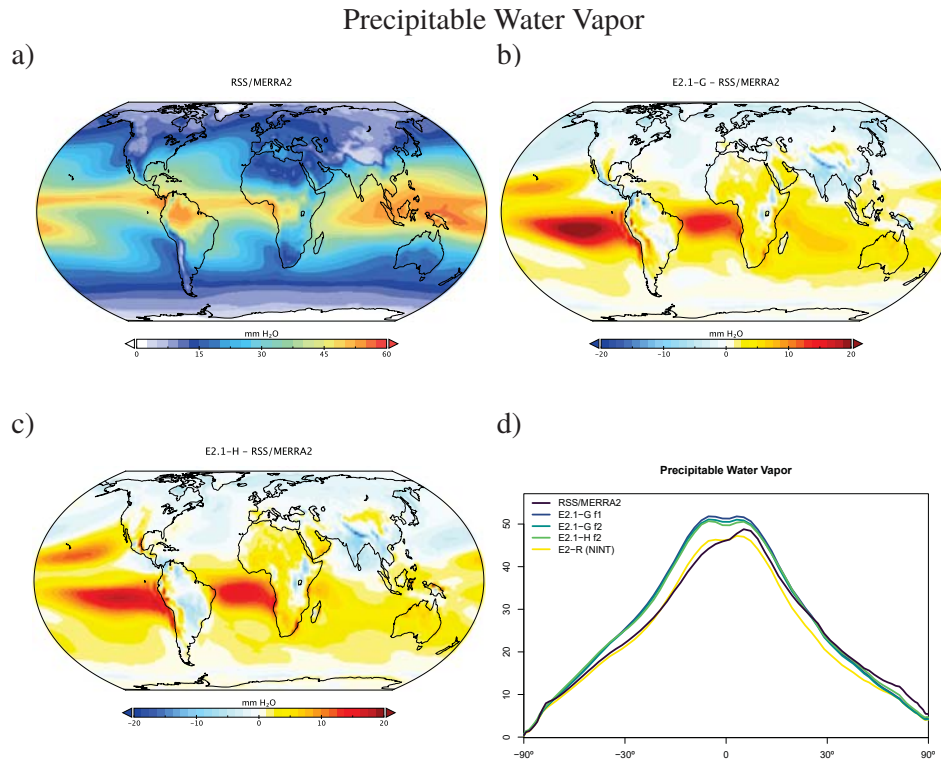
652

Figure 7. Annual climatology of short wave cloud radiative forcing, figure description as in fig. 2

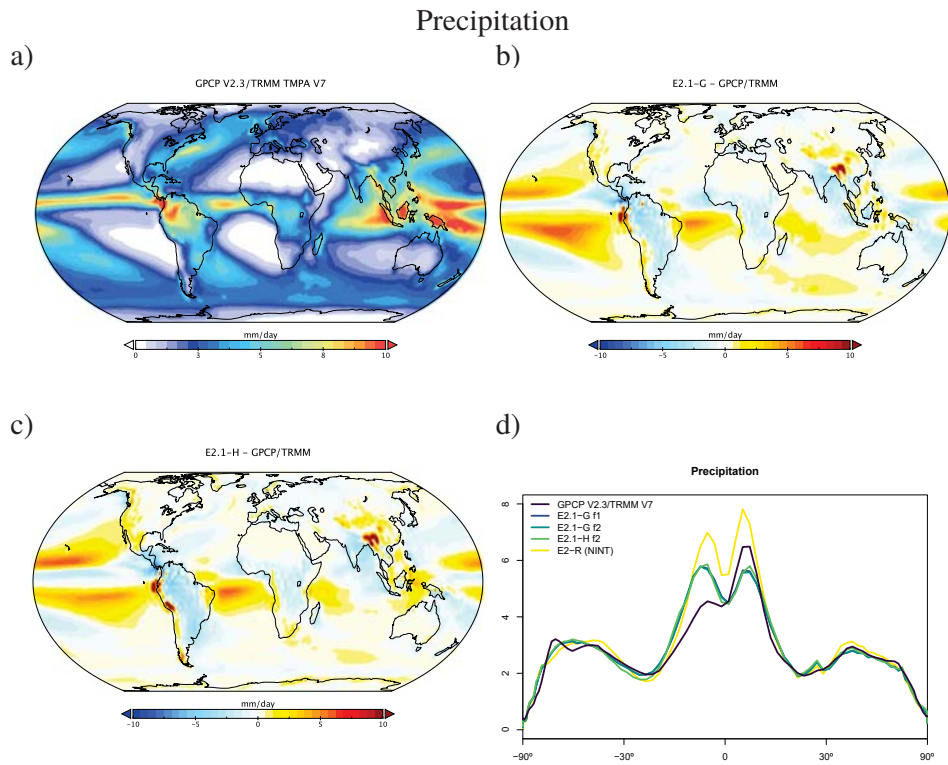


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Figure 8. Annual climatology of long wave cloud radiative forcing, figure description as in fig. 2

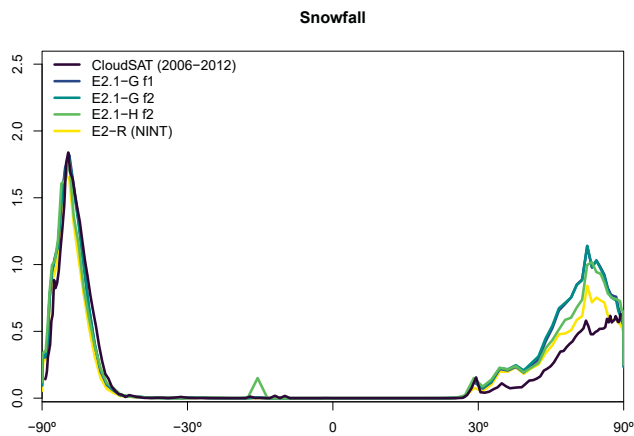


654 **Figure 9.** Annual climatology of precipitable water vapour, figure description as in fig. 2. Data derived
655 from a blend of RSS and MERRA2 products over ocean and land respectively.



656

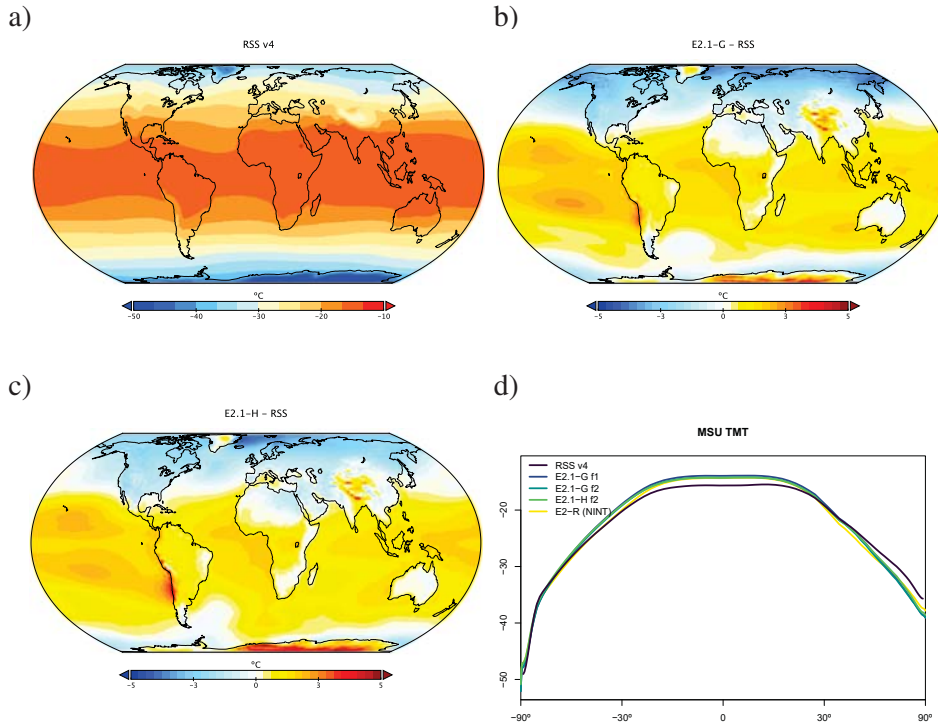
Figure 10. Annual climatology of precipitation. Figure description is as fig. 2.



657

Figure 11. Annual climatology of snowfall compared to CloudSAT data *Liu* [2008].

Mid-Tropospheric Temperature (MSU-TMT)



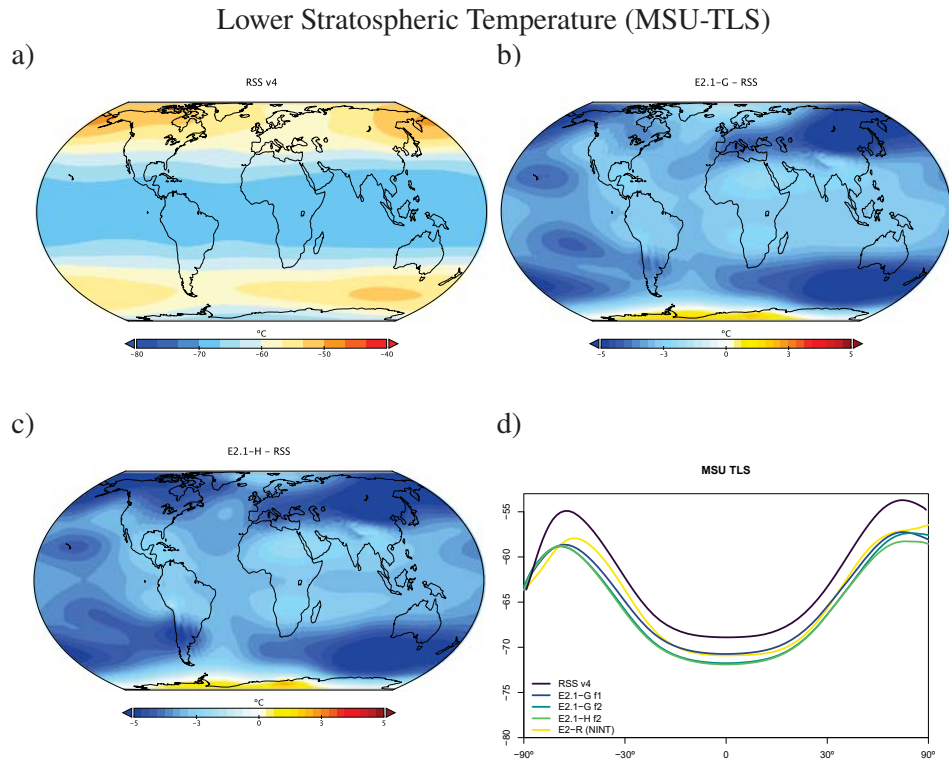
679 **Figure 12.** Annual climatology of MSU TMT. Observational data comes from RSS (1979–2014) (version
 680 4.0) [Mears and Wentz, 2016]. Figure description is as fig. 2.

5.3 Satellite-derived Atmospheric Temperatures

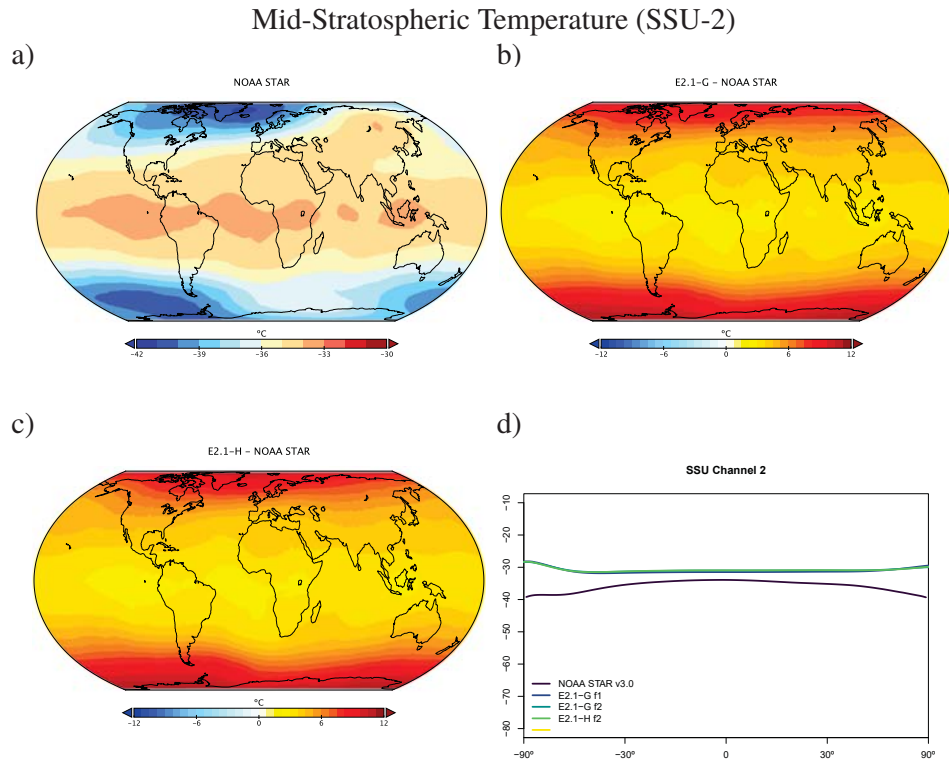
658
 659 The structure of temperature through the atmosphere plays a large role in defining
 660 fingerprints of climate change forcings, and so we compare the models to the Microwave
 661 Sounding Unit (MSU) and Stratospheric Sounding Unit (SSU) 1979–2014 brightness tem-
 662 perature climatologies (figs. 12, 13, 14). We highlight results from the mid-troposphere
 663 (TMT), the lower stratosphere (TLS) and middle stratosphere (SSU Channel 2) which have
 664 global weightings centered on 600, 70 and 4 hPa, respectively (though with substantial tails)
 665 [Mears and Wentz, 2016; Zou and Qian, 2016]. We use a static weighting function to esti-
 666 mate the channels, which though slightly less accurate than a radiative transfer calculation
 667 that takes into account surface emissivity, atmospheric water vapor, and temperature profiles
 668 [Shah and Rind, 1995], does not produce significantly different results.

669 Starting with MSU-TMT (fig. 12), the models show a notable warm bias in the trop-
 670 ics and sub-tropics, indicating a slightly steeper lapse rate in the troposphere than observed,
 671 and a cold bias in the northern high latitudes. Warm biases over high topography may be an
 672 artifact of the diagnostic comparison.

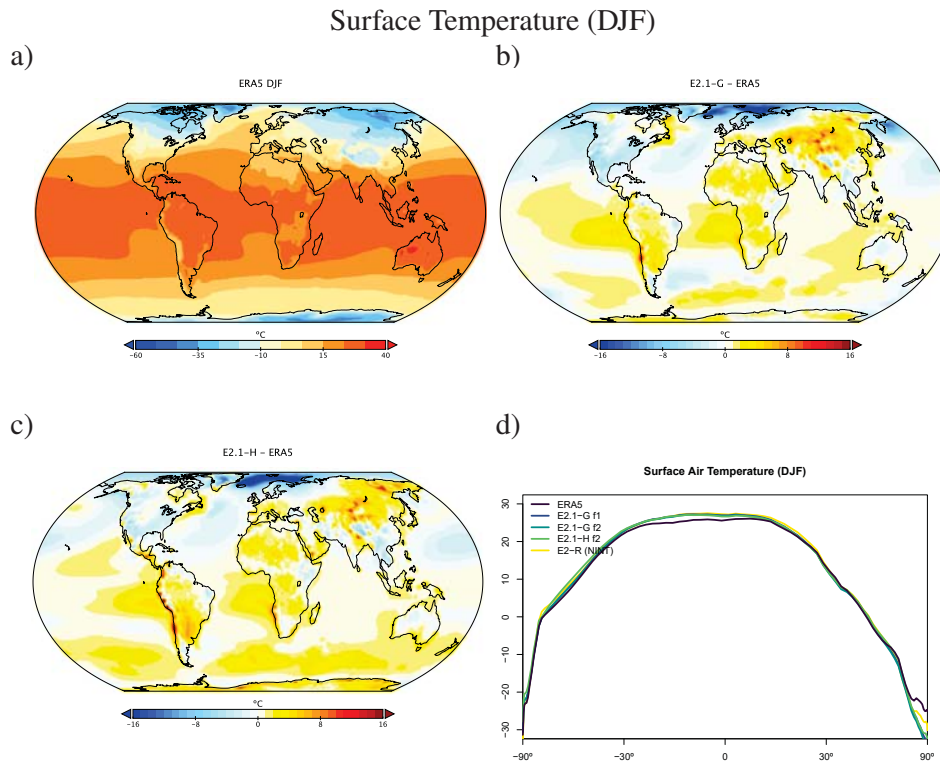
673 In the lower stratosphere (fig. 13), the models are anomalously cold, though partially
 674 the worse comparison to observations since E2 is related to a warmer climatology in the lat-
 675 est RSS version 4.0 [Mears and Wentz, 2016]. The mid and upper stratosphere (as illustrated
 676 by the SSU-2 channel, fig. 14) is too warm - particularly in the polar regions. This overall
 677 pattern of stratospheric temperature error is consistent with, but not completely explained by,
 678 a too weak Brewer-Dobson circulation in this relatively low-top model.



681 **Figure 13.** Annual climatology of MSU TLS. Observational data comes from RSS (1979–2014) (version
 682 4.0) [Mears and Wentz, 2016]. Figure description is as fig. 2.



683 **Figure 14.** Annual climatology of SSU Channel 2. Observational data comes from NOAA STAR (1979–
 684 2014) (version 3.0) [Zou and Qian, 2016]. Figure description is as fig. 2.



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Figure 15. DJF climatology of surface air temperature. Figure description is as fig. 2.

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5.4 Surface Fields

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Surface field observations are taken from ERA5 [Copernicus Climate Change Service (C3S), 2017]. Overall biases in the surface temperature fields (figs. 16 and 15) are similar to the situation seen in CMIP5, though the magnitude of errors in the Southern Oceans are notably reduced (consistent with the improvements of cloud and radiation diagnostics discussed above). Land errors are reduced, though the winter cool bias in the Arctic is slightly increased.

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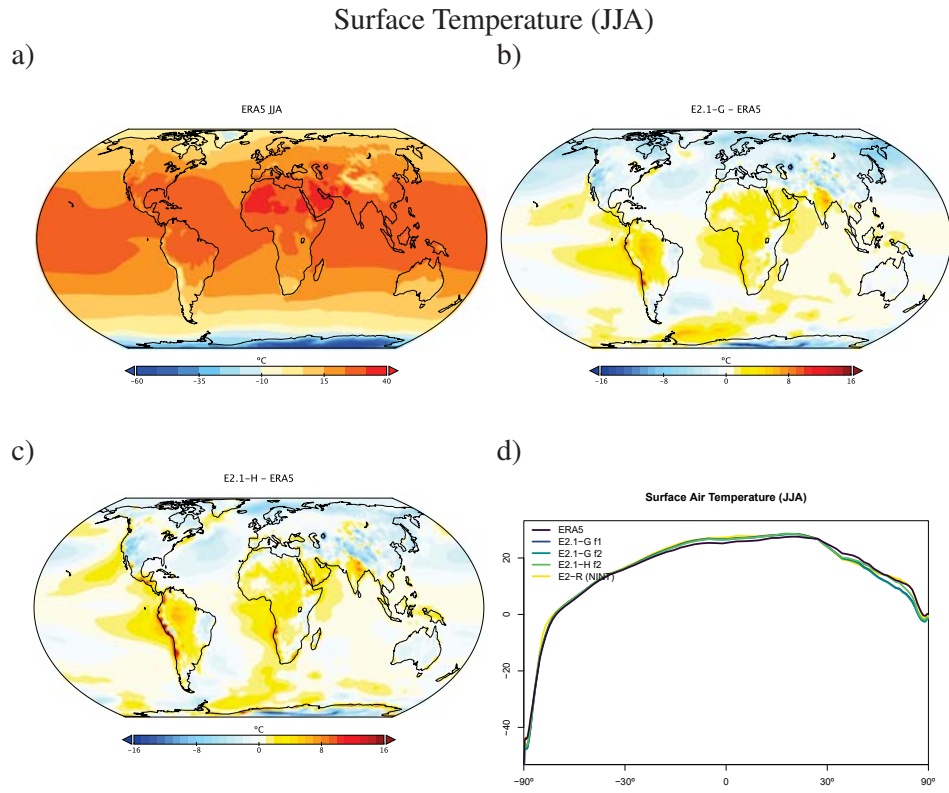
697

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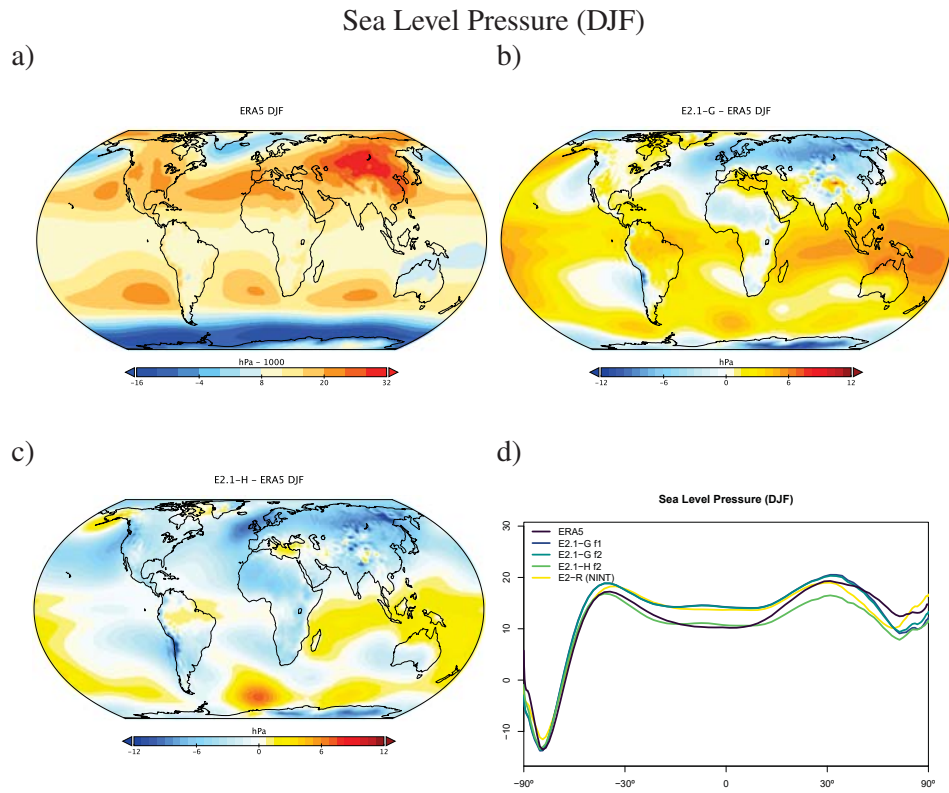
700

Sea level pressure biases are quite different between the two ocean model versions (figs. 18 and 17), with E2.1-G having more excessive high pressure in the tropics than in E2.1-H. This is partially explained by the higher than observed water vapor in the models, and the topographic change made in the HYCOM land-ocean grid which increased surface pressure over land (and through conservation, decreased it over the oceans). In the northern summer, both models fail to generate as large a jet stream gradient as observed. However, the overall pattern of surface wind stress is improved from E2 (fig. 19), with a notably more poleward and stronger maxima in the mid-to-high latitudes. There remains a westward bias in the eastern tropical Pacific.

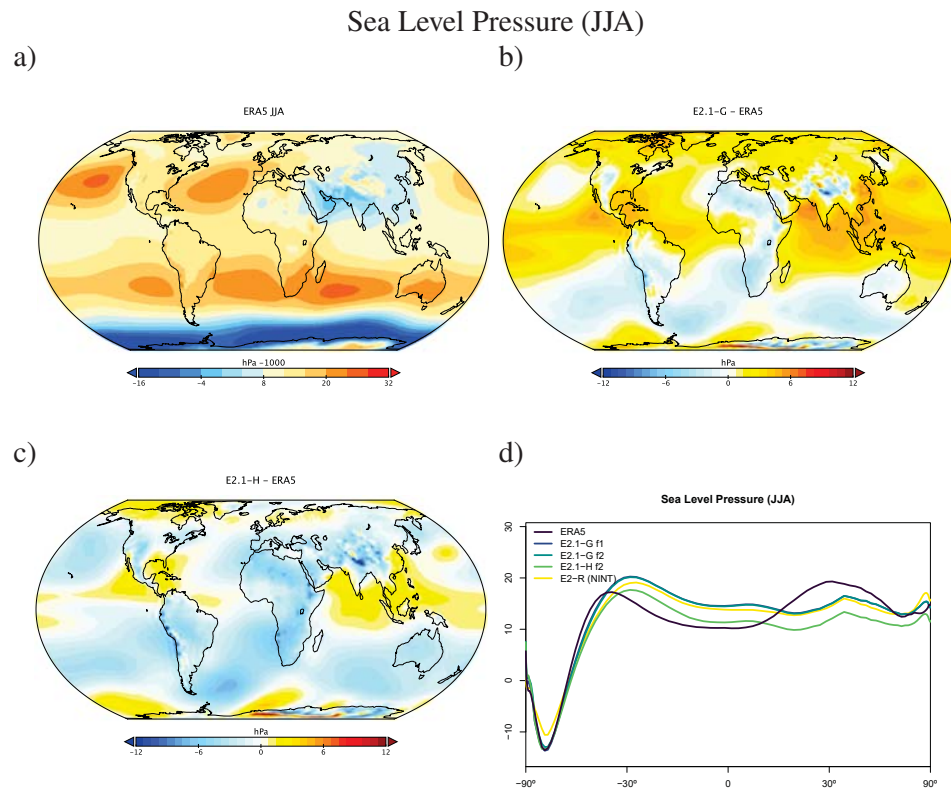


702

Figure 16. JJA climatology of surface air temperature. Figure description is as fig. 2.

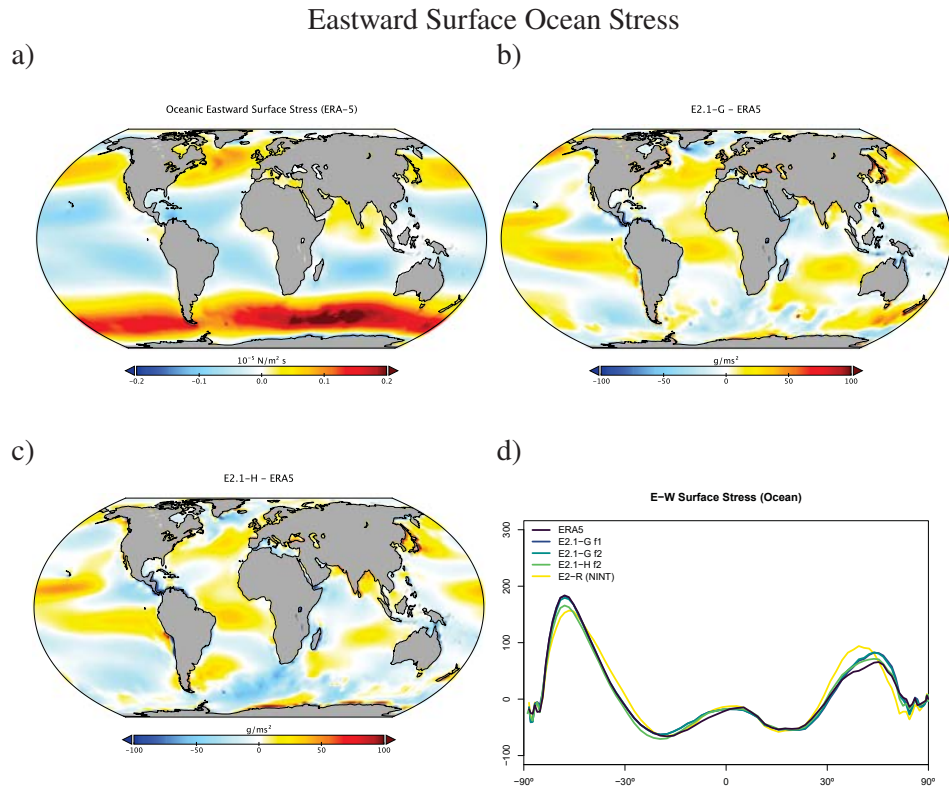


703 **Figure 17.** DJF climatology of sea level pressure (including water vapor mass in the diagnostic, even
 704 though it is not seen by the dynamics). Figure description is as fig. 2.



705

Figure 18. JJA climatology of sea level pressure. Figure description is as fig. 2.



706

Figure 19. Annual climatology of Eastward surface ocean stress. Figure description is as fig. 2.

River	E2.1-G	E2.1-H	E2-R	E2-H	Observations
Amazon	251±8	280	198–236	229–300	545
Congo	22±1	36	35–69	41–82	106
Brahmaputra-Ganges	125±4	81	68–86	110–140	105
Yangtze	108±2	111	85–100	191–210	78
Lena	44±1	41	32–34	29–31	40
Ob	52±1	38	47–52	80–89	33
St. Lawrence	56±1	35	53–55	27–28	29
Mackenzie	23±1	29	28–29	31	24

Table 5. Annual mean discharge from selected rivers ($\text{km}^3 \text{ month}^{-1}$). Ensemble mean and standard deviation for E2.1-G, range of model values from E2-R/H. Observations from *Fekete et al.* [2001].

Runoff from the major rivers can be compared to observational data [*Fekete et al.*, 2001] (Table 5). In the tropics, runoff is severely deficient in the Amazon basin and African rain forests (due to insufficient rainfall) where skill has not increased compared to earlier model versions. High latitude rivers are, however, more consistently modeled.

5.5 Ocean

We focus here on the diagnostics that most impact the coupled simulation and are straightforwardly comparable to observations. More detailed description and analysis of E2.1 ocean circulation and structure will be presented elsewhere.

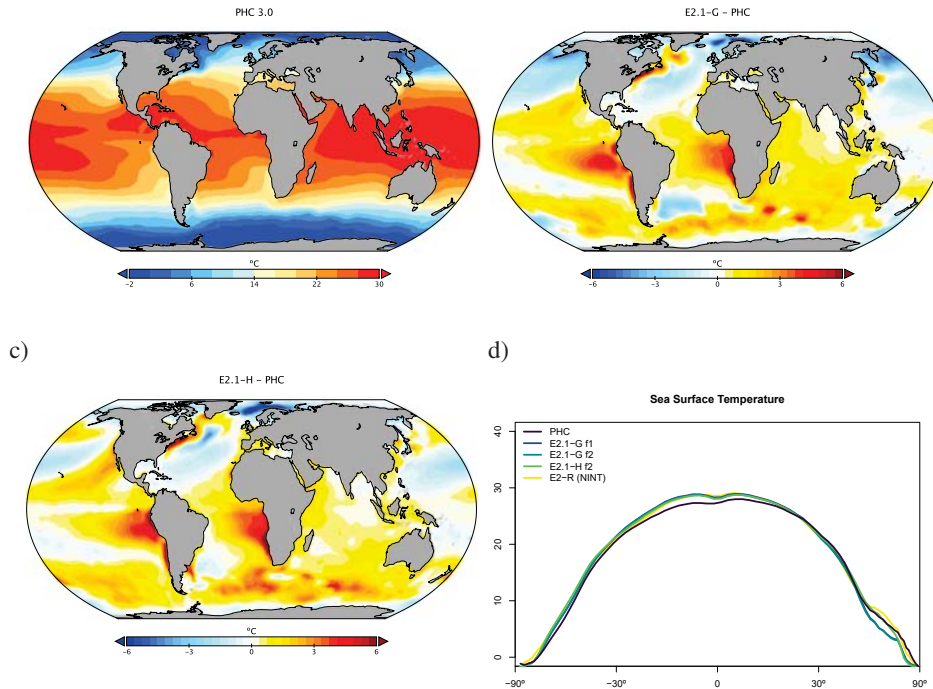
Sea surface temperature biases (fig. 20) are still dominated by the errors in the marine stratus regions, though are (again) improved in the Southern Ocean region. Arctic biases are colder than before. Overall, tropical temperatures are slightly warm, particularly in the southern tropics, which is consistent with the errors in precipitable water vapour, clouds and radiation seen above. Remarkably, the difference seen across the ocean models is quite small.

Salinity biases in E2.1-G are far smaller than in E2-R, particularly in marginal seas, but also in the open ocean (fig. 21). Clear positive biases are obvious near major river mouths (consistent with insufficient river outflow seen in Table 5).

For HYCOM, the biases in surface salinity (fig. 21c) have been totally reversed, in part due to the correction to virtual salt fluxes, from a large excess salinity in E2-H, to an overall underestimated salinity in E2.1-H, though with a reduced overall error. Arctic biases are noticeably reduced, possibly associated with the implementation of the BP ice thermodynamics.

Ocean transports are also greatly improved, notably the Drakes Passage where offsets to the observed transport are much less than previously in both models (Table 6). Fluxes through the Gulf Stream and Kuroshio Current are reasonable, but slightly higher than inferred from observations. The mass and heat transports at 26°N from the N. Atlantic overturning circulation in E2.1-H are in good agreement with direct observations [*McCarthy et al.*, 2015; *Smeed et al.*, 2019; *Johns et al.*, 2011], but larger in E2.1-G.. Poleward heat transports peak above 1 PW at $\sim 20^\circ\text{N}$, this is significantly higher than the estimates derived from a ocean state estimation approach [*Forget and Ferreira*, 2019] (fig. 22), but in reasonable agreement with direct heat flux estimates [*Ganachaud and Wunsch*, 2003]. Poleward transports in the southern oceans in E2.1-G are much more consistent with both direct measurements and ocean state estimates.

Sensitivity experiments with a reduced tidal mixing efficiency in E2.1-G suggested that tuning of this parameter could match the target Atlantic overturning transport metric at 26°N and the *Forget and Ferreira* [2019] heat transport there, but with the penalty of unacceptably



749 **Figure 20.** Annual climatology of sea surface temperature compared to the PHC 3.0 product. Figure
 750 description is as fig. 2.

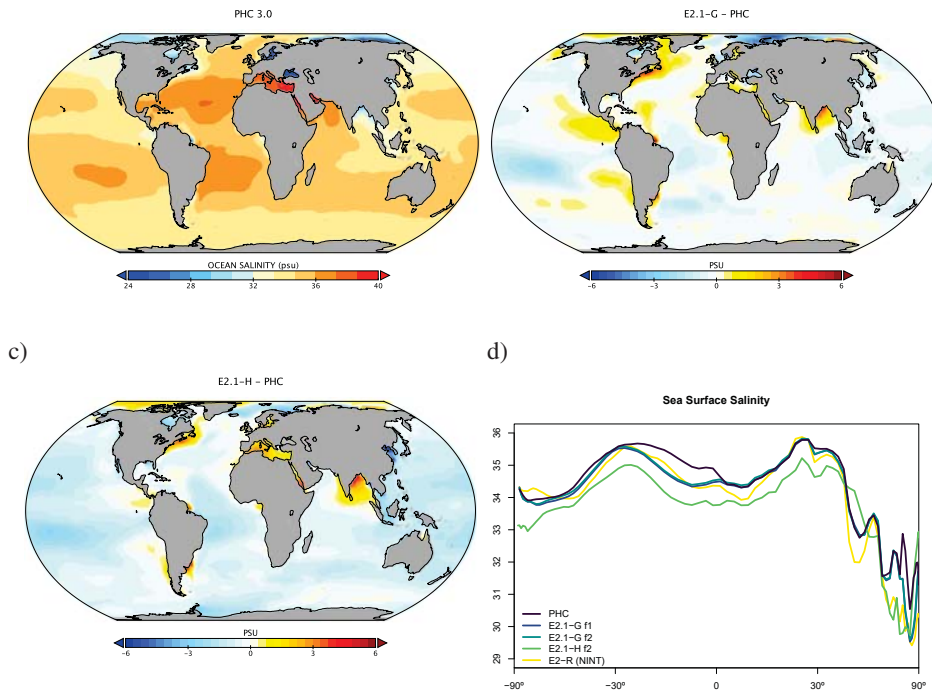
744 increasing cold biases in northern midlatitudes and the Arctic. Such compromises will be
 745 revisited in future model versions having improved cloud radiative forcing and atmospheric
 746 transports. Ocean-only experiments with an E2.1-G prototype [Romanou *et al.*, 2017] indicate
 747 that its CFC uptake is best matched in configurations having weaker AMOC magnitudes
 748 than those realized here, with implications for heat and carbon uptake.

757 **5.6 Cryosphere**

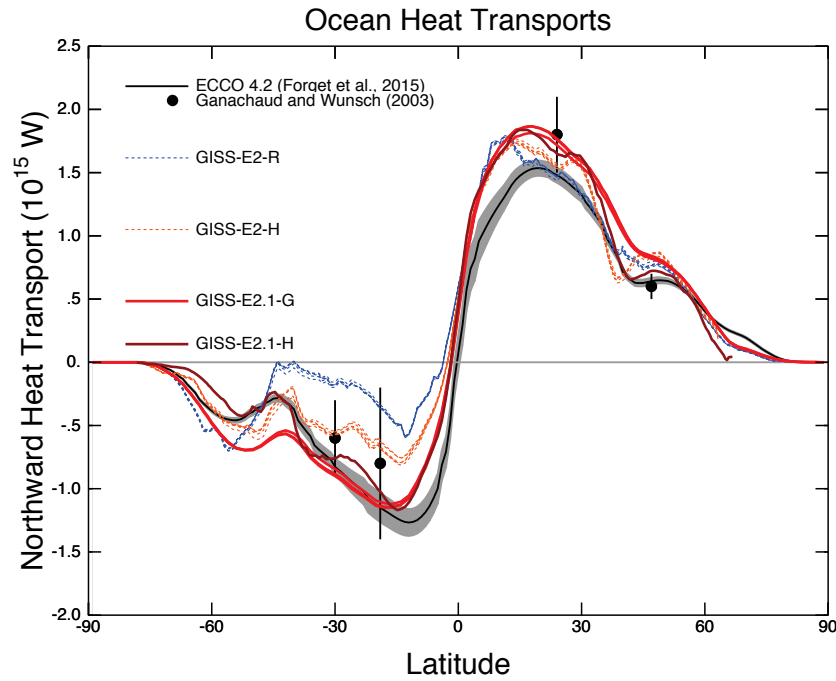
758 Figures 23 show that modeled seasonal cycles of sea ice in the Arctic and Antarctic
 759 have improved, although the absolute value of Arctic sea ice is still too large. Improvements
 760 in snow depth, meltpond extent, and albedo (as compared to SHEBA data [Curry *et al.*,
 761 2001]) are clear (fig. 24) and due to model fixes and the inclusion of the Brine-Pocket ther-
 762 modynamics scheme. Southern Ocean improvements are largely due to a more stratified
 763 ocean and an associated reduction of vertical mixing in both E2.1-G and E2.1-H, as opposed
 764 to changes in sea ice physics or properties (as has been the case previously [Liu *et al.*, 2003]).
 765 Brighter clouds in the Southern Ocean in E2.1 also cool temperatures and aid ice formation
 766 there.

Diagnostic	E2.1-G	E2.1-H	E2-R	E2-H	Observations
N. Atl. MOC (Max)	27.2	20.4±0.3	27.2±0.7	24.5± 0.8	-
N. Atl. MOC (26°N)	24.8±0.4	17.8±0.3	18.4±0.3	22.4±0.6	≈18 ^{R19}
Atl. Heat (26°N)	1.21±0.01	1.06±0.01	0.97±0.01	0.99±0.02	1.3±0.4 ^{J11} /0.88±0.01 ^{F19}
ACC (Drake Pass.)	150±1	178±1	254± 1	192±2	130 ^{P88} /173 ^{D16}
Gulf Stream	55±1	48.2±0.3	49± 1	39.8±0.8	≈35 ^{R11}
Kuroshio	49±1	67±2	64±1	71.7±0.5	≈57 ^{I01}
Bering Str.	0.16±0.003	-0.55±0.01	0.16±0.01	0.45±0.01	0.8±0.2 ^{W05}
Indonesian throughflow	18.9±0.2	18.4±0.2	11.5±0.2	17.6±0.3	15 ^{S09}

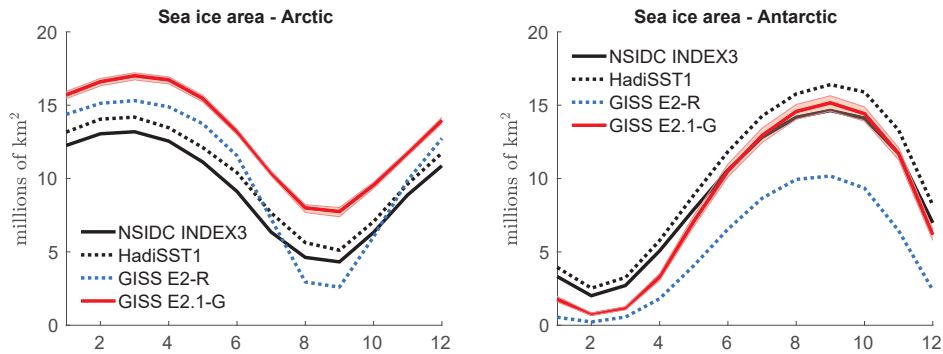
Table 6. Selected ocean mass (Sv) and heat (PW) fluxes. Range is standard deviation of the 1979–2014 average from 5 ensemble members for each configuration. Observations: ^{R19} McCarthy et al. [2015]; ^{Smeed} Smeed et al. [2019] (estimate at 26°N); ^{P88} Petersen [1988]; ^{D16} Donohue et al. [2016]; ^{J11} Johns et al. [2011]; ^{I01} Imawaki et al. [2001]; ^{W05} Woodgate et al. [2005]; ^{S09} Sprintall et al. [2009]; ^{F19} Forget and Ferreira [2019]



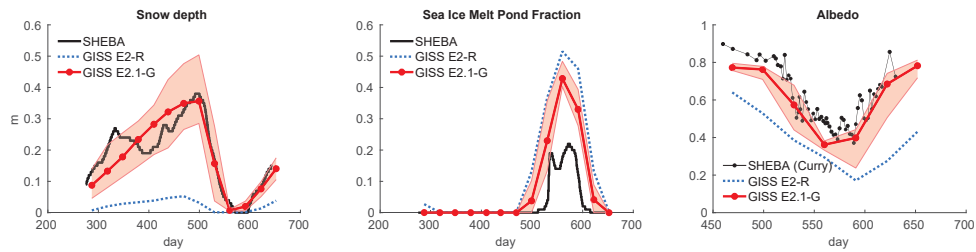
751 **Figure 21.** Annual climatology of sea surface salinity compared to the PHC 3.0 product. Figure description
752 is as fig. 2.



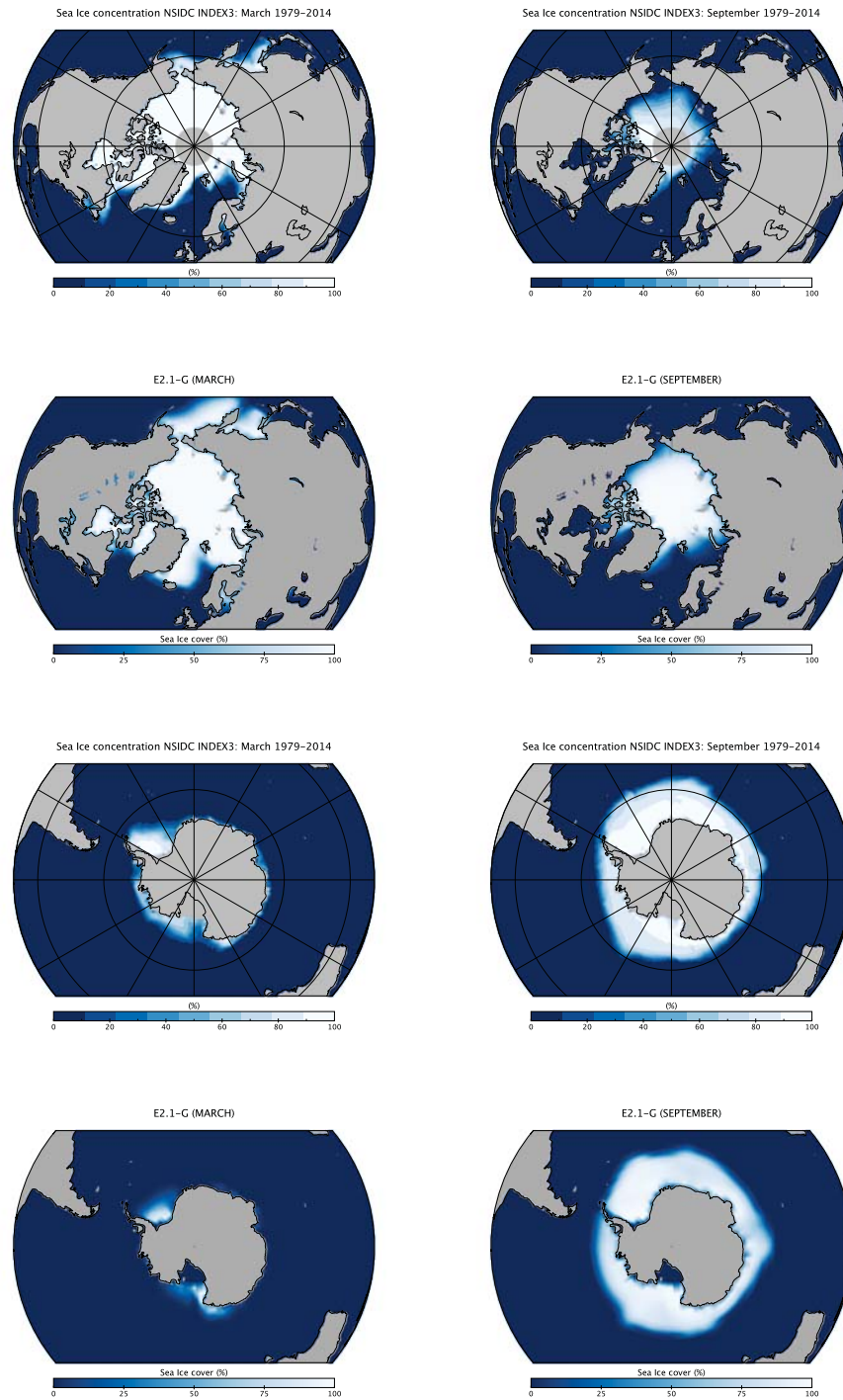
753 **Figure 22.** Annual mean global northward ocean heat transports. Comparisons of the models with mean
 754 estimates from 1992–2011 from the ECCO ocean state estimate (v4 release 2) with 95% confidence intervals
 755 on the mean derived from the interannual variability [Forget *et al.*, 2015; Forget and Ferreira, 2019] and
 756 oceanographic estimates [Ganachaud and Wunsch, 2003].



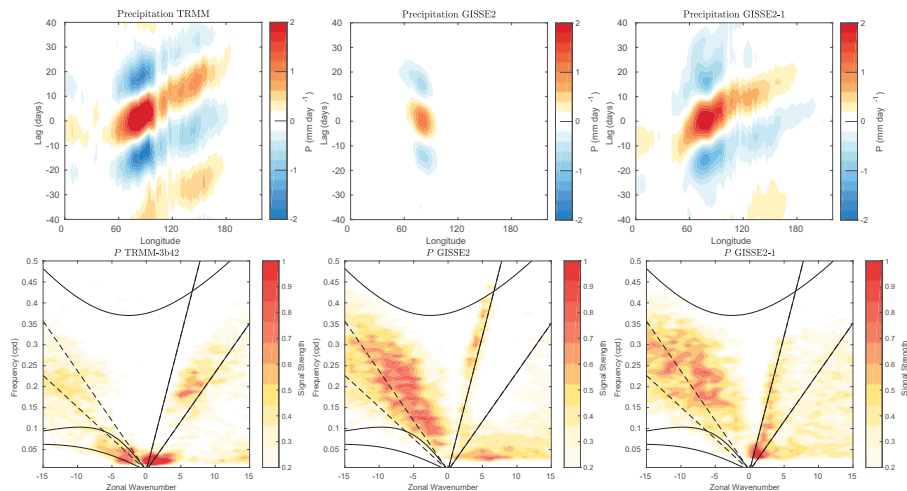
767 **Figure 23.** Annual climatology of sea ice area in both hemispheres in E2-R (blue dashed) and E2.1-G (red).
 768 Observational data comes from NSIDC (1979–2014), after correction for the Arctic polar 'hole' [Fetterer
 769 *et al.*, 2011] and HadISST1 (1979–2014) [Rayner *et al.*, 2011]. The ensemble mean climatology is plotted for
 770 E2-R (1979–2012) and E2.1-G (1979–2014), with spread across E2.1-G ensemble members in pink.



771 **Figure 24.** Spot comparisons of the E2-R (blue dashed) and E2.1-G (red) simulations against the SHEBA
 772 measurements for snow depth, meltpond fraction and albedo [Curry *et al.*, 2001]. Ensemble spread for E2.1-
 773 G is in pink.



774 **Figure 25.** Sea ice concentration (%) for March (left column) and September (right column) in the NSIDC
 775 observations and E2.1-G simulations. Rows a) and b) Arctic, and rows c) and d) Antarctic. E2.1-H results are
 776 similar.



788 **Figure 26.** Comparison of MJO signals and propagation in the TRMM data (release 3b24), [Iguchi
 789 *et al.*, 2000] and in E2-R and in E2.1-G simulations. (Top) Hovmöller plots of MJO propagation. (Bot-
 790 tom) Wheeler-Kiladis diagrams for tropical wave motion [Wheeler and Kiladis, 1999]. Figures courtesy of
 791 Ángel Adames.

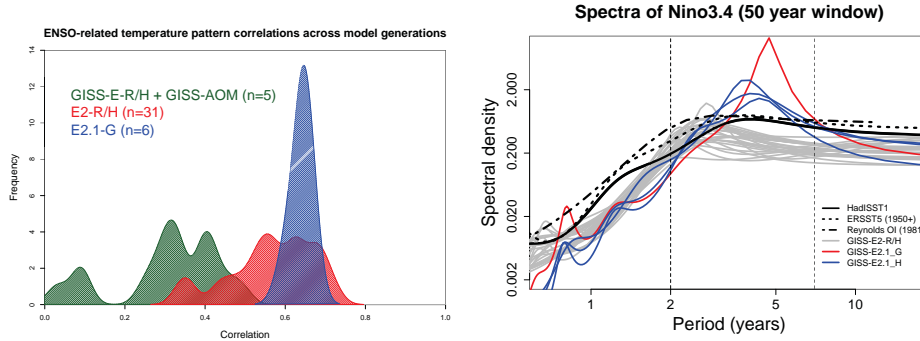
777 **5.7 Model internal variability**

778 As model processes have become more sophisticated and the base climatology has be-
 779 come more realistic, the realism of the patterns of internal variability has also improved. We
 780 focus here on ENSO, the PDO and the MJO because the improvements over previous mod-
 781 els have been most dramatic. Notably, while the MJO was a specific target for model im-
 782 provement through the model development process, the changes in ENSO and PDO patterns
 783 emerged as part of the overall improvement in skill.

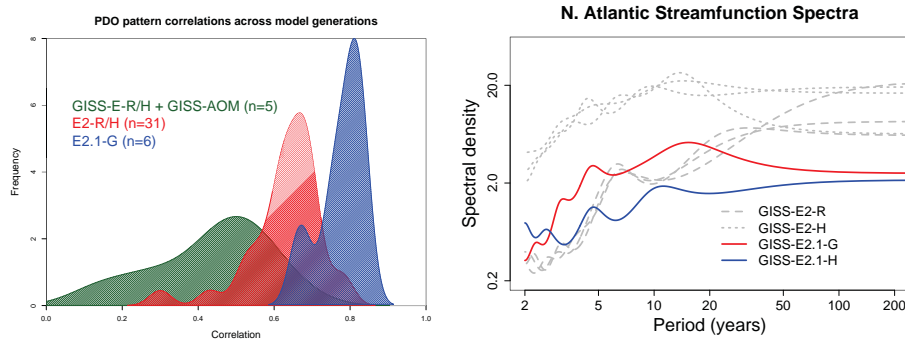
784 The MJO improvement is highlighted in figure 26, where the lack of MJO-related activ-
 785 ity and lack of propagating features in the Pacific in E2-R was very clear in comparison
 786 with an analysis of the TRMM data. However, in E2.1-G, the improvement in propagation
 787 and in the wavenumber/frequency plot [Wheeler and Kiladis, 1999] is evident.

800 For the longer term tropical modes, ENSO and the PDO, there have been large im-
 801 provements in the patterns of associated temperature variability (figs. 27a and 28a) across
 802 CMIP generations, and particularly since CMIP5. However, that improvement must be tem-
 803 pered by a recognition that the spectral signature of ENSO has not improved. In all versions
 804 of E2, there was insufficient overall variance, and in particularly a deficit in variability at 3–
 805 7 years (overall standard deviations were 0.60°C for E2-R and 0.67°C for E2-H, compared
 806 to ~1°C in the ERSST5 observations). However, in E2.1-G and E2.1-H the 2 to 4-year vari-
 807 ability is now too strong (figs. 27b). The overall Nino3.4 standard deviation is modestly too
 808 high (1.2°C) in E2.1-G and too low in E2.1-H (0.75°C), compared to ~1°C in the ERSST5
 809 observations [Huang *et al.*, 2017]. This nonetheless impacts the interannual variability in
 810 these simulations, even for the global mean, leading us to increase the number of ensemble
 811 members in the E2.1-G historical simulations in order to better assess the forced responses.

812 The larger overall ENSO variability and unrealistically peaked spectral signature in
 813 E2.1-G relative to E2.1-H suggest that ocean interior structure and damping mechanisms
 814 exert as much influence as atmospheric processes. Some of the latter have been quantified
 815 in feedback form for E2.1-G following Figure 7 in Bellenger *et al.* [2014]. Specifically, the
 816 wind-stress (positive) feedback is $9.8 \times 10^{-3} \text{ N m}^{-2} \text{ } ^\circ\text{C}^{-1}$, 20% weaker than in ERA40, and



792 **Figure 27.** a) Improvement of pattern correlations of ENSO to the observations over GISS model gener-
 793 ation (from CMIP3 to CMIP6). Calculations via the Climate Variability Data Portal (CVDP [Phillips *et al.*,
 794 2014]). b) Spectra of Nino3.4 variability in 50 year segments from the PI-controls compared to various
 795 observational products.



796 **Figure 28.** a) Improvement of pattern correlations of the PDO to the observations over GISS model gener-
 797 ation (from CMIP3 to CMIP6). Calculations via the Climate Variability Data Portal (CVDP [Phillips *et al.*,
 798 2014]). b) Spectra of variability in the N. Atlantic annual mean maximum streamfunction (derived from a
 799 detrended 1000 years of PI-control simulation).

817 the surface-flux (negative) feedback is $-12.5 \text{ W m}^{-2} \text{ }^{\circ}\text{C}^{-1}$, 30% weaker than observed. In a
 818 sensitivity test (similar to one reported in Rind *et al.* [2020]), we applied a change to the at-
 819 mospheric convection scheme that led to reduced ENSO amplitude and a shift of the peak
 820 to shorter periods. Both of the feedback coefficients are significantly smaller in that simu-
 821 lation, suggesting that its ENSO improvement occurred for the wrong reasons, and overall
 822 model skill was not enhanced. This remains an active area of model testing, although we an-
 823 ticipate that it will require a substantial improvement of marine stratus biases (as a function
 824 of increased vertical resolution and better moist physics) before tuning for the correct ENSO
 825 feedbacks will become worthwhile.

826 In the North Atlantic, where decadal and greater variability is associated with the over-
 827 turning streamfunction there are mixed changes. There is greater variability at 8–15 yrs for
 828 E2.1-G compared to E2-R, but significantly less variability in E2.1-H compared to E2-H.
 829 The standard deviation of the detrended annual streamfunction maximum is 1.4 Sv for E2.1-
 830 G, and 1.0 for E2.1-H (compared to 1.3 and 3.2 Sv in E2-R and E2-H respectively). This can

831 be compared to the interannual variability in the meridional overturning circulation at 26°N
832 [*McCarthy et al.*, 2015; *Smeed et al.*, 2019] of ~1.3 Sv.

Field	E2.1-G (f2)	E2.1-G (f1)	E2.1-H (f2)	E2-R	E2-H
OTR	0.68	0.68	0.67	0.66	0.63
ASR	0.84	0.84	0.85	0.79	0.78
MSU-TMT	0.89	0.88	0.90	0.90	0.90
MSU-TLS	0.64	0.69	0.62	0.73	0.71
TOTAL CLOUD	0.32	0.32	0.31	0.19	0.17
LOW CLOUD	0.35	0.35	0.34	0.16	0.12
SLP (DJF)	0.76	0.75	0.81	0.78	0.71
SLP (JJA)	0.82	0.82	0.83	0.79	0.75
SAT (DJF)	0.90	0.90	0.89	0.90	0.88
SAT (JJA)	0.90	0.90	0.90	0.90	0.87
PRECIP	0.52	0.51	0.51	0.50	0.45
EWSS	0.81	0.81	0.77	0.78	0.71
SST	0.90	0.90	0.90	0.91	0.86
SSS	0.73	0.73	0.57	0.63	0.54

857 **Table 7.** Arcsin-Mielke scores across model configurations for selected fields as referenced above (see
858 fig. 29 for the field definitions, though we also add sea surface temperature (SST) and salinity (SSS) in this
859 table). The highest score across the coupled models for each field is highlighted. Note that for the E2 models,
860 the output data is from 1979–2004, while the target climatologies are as described above.

833 5.8 Summary Statistics

834 We are interested both in how model evolution affects skill scores, but also in how the
835 GISS model compares to similarly functional models in the CMIP5 and CMIP6 ensembles.
836 Improvements across the board are seen in standard large scale climatological tests as seen
837 in the Taylor diagram comparing E2-R (NINT) with E2.1-G (NINT) (Fig. 29). The improve-
838 ments are largest in fields that were the worst performing in CMIP5 (clouds, precipitation),
839 though still positive for even well-simulated fields. As in previous papers, we can calculate
840 an Arcsin-Mielke score (between 0 and 1) [Watterson, 1996] for a suite a standard variables
841 (Table 7). These reflect the same general tendencies. Differences between the f1 and f2 en-
842 sembles are barely perceptible (except for MSU-TLS which is a little better in the f1 ensem-
843 ble).

844 Any overall ranking of performance is by necessity ad hoc, and not determinative of
845 every metric, but across a range of measures, the E2.1-G (f2) is the best performing con-
846 figuration considered here. There are small degradations of skill for the MSU diagnostics
847 (though not for the trends [Miller et al., 2020]). E2.1-H has slightly better SLP patterns, but
848 the differences in atmospheric variables is minor, especially compared to the improvements
849 of all E2.1 configurations with respect to E2.

861 6 Climate Sensitivities

862 As part of the DECK simulations requested by CMIP6, we performed a number of ide-
863 alized simulations (1pct4xCO2, abrupt4xCO2) as well as some related simulations (abrupt2xCO2
864 with the coupled and q-flux ocean versions) (all performed with the f1 background compo-
865 sition). The summary of various metrics of climate sensitivity (along with the comparison
866 to the previous models) is seen in Table 8. We note that the effective CS as calculated by the
867 Gregory method [Gregory et al., 2004] almost always underestimates the true long term ECS
868 by 10 to 20% [Dunne et al., 2020]. The perhaps more relevant TCR is slightly larger in the
869 E2.1 models than previously, consistent with a smaller rate of mixing of heat into the ocean
870 (and slightly smaller overall radiative imbalance (Table 2)

Model version & configuration	ECS _{qflux}	CS _{Eff}	ECS		TCR
			from 4×CO ₂	from 2×CO ₂	
E2.1-G (NINT)	3.0	2.7	3.2	3.6	1.8
E2.1-H (NINT)	"	3.1	3.5	3.4	1.9
E2.1-G (OMA)	2.9	2.6			1.6
E2.1-H (OMA)	"	3.1			2.0
E2.1-G (MATRIX)					1.8
E2.1-G (MATRIX)					2.0
E2.1-G (TOMAS)	3.1				
E2-R (NINT)	2.7	2.1	2.3	2.6	1.4
E2-H (NINT)	"	2.3	2.5		1.7
E2-R (TCADI/OMA)	3.0	2.4			1.6
E2-H (TCADI/OMA)	"	2.5			1.8

871 **Table 8.** Climate Sensitivities to 2×CO₂ (°C) estimated multiple ways (note that not all calculations have
872 been completed with all versions). Equilibrium Climate Sensitivity (ECS) is defined from multi-millennial
873 coupled simulations, or from a q-flux (slab ocean) model (ECS_{qflux}). CS_{Eff} is from a linear extrapolation of
874 yr 1–150 results in the abrupt4xCO₂ simulations [Gregory *et al.*, 2004]. Transient Climate Responses (TCRs)
875 are taken from year 70 in the 1pct4xCO₂ simulation.

876 The relative stability of the climate sensitivity from E2 to E2.1 is however due to two
877 counteracting influences. First, as discussed in Miller *et al.* [2020], the effective radiative
878 forcing associated with a doubling of CO₂ is 15% smaller (3.59 compared to 4.19 W m⁻²) in
879 the E2.1 model than it was in E2 and closer to the canonical 3.7 W m⁻² [Myhre *et al.*, 2013].
880 This is mostly explained by higher water vapor content and greater LW cloud forcing which
881 reduce the baseline contribution of CO₂ to longwave opacity, and hence reduce the sensitiv-
882 ity to opacity changes. Secondly, the changes to cloud feedbacks associated with the increase
883 in supercooled cloud water make the overall cloud feedbacks more positive (by reducing the
884 negative cloud phase feedback). Thus the impact to 2×CO₂ is only slightly changed, though
885 the normalised sensitivity has increased substantially from 0.62 to 1.00 °C W⁻¹ m² (using
886 the ECS from 2×CO₂), or similarly from 0.58 to 0.87 W⁻¹ m² (using the long-term response
887 to 4×CO₂).

888 7 Conclusions

889 As computational resources increase, the temptation at many climate modeling centers
890 is to increase resolution (and therefore compute time) such that the overall throughput of the
891 model stays roughly constant. In contrast to that strategy though, the increment from GISS-
892 E2 to GISS-E2.1 versions focused instead on fixes, better calibrations and in a few cases,
893 improved parameterizations. This was embarked on in parallel with a far more extensive up-
894 grade to the E3 code (including, new topologies, new dynamical cores, higher horizontal and
895 vertical resolution, and new moist physics) which will be reported elsewhere. The question
896 then arises, as to whether the first strategy can provide a worthwhile increase in skill with
897 negligible costs of additional runtime, more efficiently than the second. The answer to that is
898 a definitive yes.

899 Skill scores in E2.1 are consistently (though not universally) higher in fields that were
900 specifically tuned for as well as in emergent properties (such as the PDO patterns) that were
901 not. Improvements are physically coherent across fields, particularly in the Southern Ocean
902 where the most change has been seen in the ocean, atmosphere and cryosphere.

903 Nonetheless, we also note the limitations of this approach and the stubborn persistence
 904 of long-term biases. Notably, while cloud properties improved, the lack of sufficient marine
 905 stratus is still apparent. Similarly, the persistence of a double ITCZ, and excessive symme-
 906 try in the zonal mean tropical diagnostics has not been ameliorated to any significant extent.
 907 These features have however been almost eliminated in the E3 simulations which have had
 908 the benefit of higher resolution, greatly improved moist physics and more comprehensive
 909 calibration [Cesana *et al.*, 2019]. It's also apparent that minor retunings are not able to com-
 910 pensate for a model top that is too low for a realistic stratospheric circulation or to generate a
 911 quasi-biennial oscillation [Orbe *et al.*, 2020; Rind *et al.*, 2014].

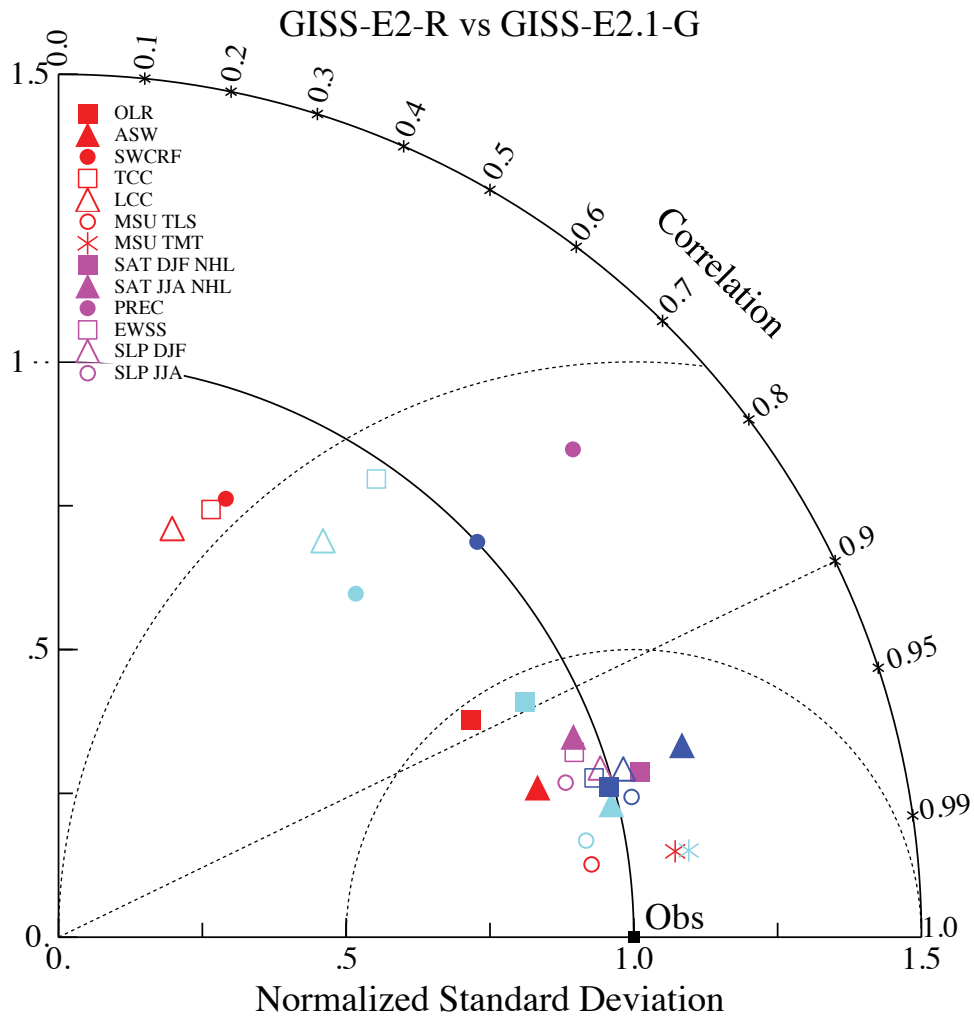
912 Within the broader constellation of the multi-model ensembles used in CMIP, true
 913 structural diversity continues to be a necessary component for any multi-model projection
 914 to have a hope of spanning the 'truth' [Knutti *et al.*, 2013]. Better-calibrated lower resolu-
 915 tion models and more sophisticated higher resolution models here can play a significant role
 916 in expanding that diversity and avoiding the potential danger of similar, and perhaps prob-
 917 lematic, new assumptions being adopted by all model groups as they jointly improve such
 918 features as cloud and aerosol microphysics [Gettelman *et al.*, 2019; Andrews *et al.*, 2019;
 919 Golaz *et al.*, 2019]. The apparent increase in climate sensitivity to doubled CO₂ in some of
 920 the next-generation models [Forster *et al.*, 2019; Dunne *et al.*, 2020] whether realistic or not,
 921 is very concerning. If this is a reflection of the real world, climate impacts are likely to be
 922 greater than we have up to now anticipated, and if it is not, then it raises serious questions
 923 about model independence and underlines the importance of true structural diversity. We
 924 simply note that the model sensitivity seen in the E2.1 models ($\sim 3^\circ\text{C}$) is near the center of
 925 the traditionally accepted range of 1.5 to 4°C. While the justification for that range has im-
 926 proved enormously since the Charney report [Charney *et al.*, 1979], the latest assessments do
 927 not challenge it [Sherwood *et al.*, 2020].

928 **8 Data and code availability**

929 All data from the piControl, historical, abrupt4xCO₂, and 1pctCO₂ simulations dis-
 930 cussed here are publicly available in the CMIP6 archive through multiple nodes of the Earth
 931 System Grid Federation. The code used corresponds to the E2.1 tag in the ModelE git reposi-
 932 tory available from the NCCS CDS system.

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 942 by the NOAA Climate and Global Change Program and are available at www.remss.com.
 943 SSU data are from NOAA / NESDIS Center for Satellite Applications and Research. Ocean
 944 Heat Content data were taken from NOAA NCEI [https://www.nodc.noaa.gov/OC5/3M_](https://www.nodc.noaa.gov/OC5/3M_HEAT_CONTENT/)
 945 HEAT_CONTENT/.



850 **Figure 29.** Summary Taylor diagram for selected quantities showing the difference in performance for
 851 E2.1-G (light and dark blue symbols) compared to E2-R (red and purple) for different fields. The change in
 852 performance in each field can be tracked by going from the red (purple) symbol to the corresponding light blue (dark blue)
 853 one. Data sources: CERES EBAF 4d1b: Outgoing Longwave Radiation (OLR) and Absorbed Solar Radiation
 854 (ASR) (60°S–60°N); RSS v4 MSU-TMT and MSU-TLS; ISCCP-H Total Cloud Cover, Low Cloud Cover
 855 (60°S–60°N), ERA-5 oceanic Sea Level Pressure (SLP) (DJF and JJA), SAT (DJF and JJA) and Eastward
 856 Surface Stress (EWSS); TRMM/GPCP Precipitation.

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